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**Near Infrared Laser Propagation and Absorption Analysis in Tissues  
Using Forward and Inverse Monte Carlo Methods**

**APPROVED BY  
SUPERVISING COMMITTEE:**

**Supervisor:**

---

Halil Berberoglu

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Yuebing Zheng

**Near Infrared Laser Propagation and Absorption Analysis in Tissues  
Using Forward and Inverse Monte Carlo Methods**

**by**

**Babak Nasouri, B.S.**

**Thesis**

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## **Dedication**

Dedicated to my parents, for their endless support and encouragement

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# **Near Infrared Laser Propagation and Absorption Analysis in Tissues Using Forward and Inverse Monte Carlo Methods**

Babak Nasouri, M.S.E.

The University of Texas at Austin, 2014

Supervisor: Halil Berberoglu

For understanding the mechanisms of low level laser/light therapy (LLLT), accurate knowledge of light interaction with tissue is necessary. In order to have a successful therapy, laser energy needs to be delivered effectively to the target location which depending on the application can be within various layers of skin or deeper. The energy deposition is controlled by input parameters such as wavelength, beam profile and laser power, which should be selected appropriately. This thesis reports a numerical study that investigates the laser penetration through the human skin and also provides a scale for selection of wavelength, beam profile and laser power for therapeutic applications.

First, human skin is modeled as a three-layer participating medium, namely epidermis, dermis, and subcutaneous, where its geometrical and optical properties were

obtained from the literature. Both refraction and reflection are taken into account at the boundaries according to Snell's law and Fresnel relations. Then, a three dimensional multi-layer reduced-variance Monte Carlo tool was implemented to simulate the laser penetration and absorption through the skin. Local profiles of light penetration and volumetric absorption densities were simulated for uniform as well as Gaussian profile beams with different spreads at 155 mW average power over the spectral range from 1000 nm to 1900 nm. The results showed that lasers within this wavelength range could be used to effectively and safely deliver energy to specific skin layers as well as to achieve large penetration depths for treating deep tissues, without causing any skin damage. In addition, by changing the beam profile from uniform to Gaussian, the local volumetric dosage could be increased as much as three times for otherwise similar lasers.

In the second part of this thesis, a three-dimensional single-layer reduced-variance inverse Monte Carlo method was developed to find the optical properties of the skin using the experimental values of transmittance and reflectance. The results showed that both transmittance and reflectance scale well with transport optical thickness. Moreover, it was also shown that penetration depth is highly sensitive to the laser wavelength and varied within the range from 1.7 mm to 4.5 mm.

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# Chapter 1

## Introduction

### 1.1 BACKGROUND AND MOTIVATION FOR THE STUDY

As an effective method for improving tissue repair, pain reduction, and hair growth, Low Level Laser Therapy (LLLT) has become a progressively dominant therapy recently. In this method, lasers with the low power range from  $10^0$  to  $10^3$  mW and the wide spectral range from 300 nm to 10,600 nm are irradiated on the tissue [1]. The success of LLLT in expediting the healing rate of skin wounds has been vastly reported in the literature for both human and animal cases [1-8]. Fillipin *et al.* [2] and Rizzi *et al.* [3] reported that, using 904 nm gallium arsenide (Ga-As) laser LED with the average low power of 45 mW, considerable reduction was achieved in oxygen release and inflammatory responses of muscle, both of which are caused by trauma. Moreover, Hopkins *et al.* [4] showed that LLLT enhanced the healing speed on human forearm abrasion using laser diodes with wavelengths between 660 nm and 940 nm and power of 15 mW to 25 mW. Similarly, a noticeable increase in the rate of cell growth in the skin of both humans [5] and pigs [6] have been observed by using low-level lasers.

Despite the success of LLLT, the numbers of ineffective cases are not negligible. For example, da Rosa *et al.* [9] reported that an 808 nm As-Ga-Al laser (100 mW, 3.57 W/cm<sup>2</sup>, 40 s) enhanced the healing process of osteoarthritis more effectively than a

660nm In-Ga-Al-P laser (100 mW, 3.57 W/cm<sup>2</sup>, 40 s), based on the experiments on 36 male adult Wistar rats. On the other hand, Taradaj *et al.* [10] claimed that LLLT application does not improve healing of venous leg ulcers in surgically and conservatively treated patients, according to their experiments with the 810 nm Ga-Al-As laser (4 J/cm<sup>2</sup>, 65 mW) on 83 patients. These results indicate that using LLLT in a “shot in the dark” manner cannot make this method a reliable therapy method. The dependence of laser power absorption and penetration depths on input parameters, such as wavelength, power and beam profile, should be investigated more precisely to improve the understanding of laser-tissue interaction. As an approach to resolve this issue, laser penetration through tissues was studied numerically using single-layer model for the human skin [11-19]. For example, Roeva *et al.* [17] investigated laser penetration by applying LLLT on 142 patients and also evaluating the light delivery within the tissue using Monte Carlo method. They simulated the human tissue as a single layer medium with the constant optical properties, and showed the effect of absorption and scattering coefficients on the laser penetration in skin. Moreover, Stoykova *et al.* [18] considered a single layer model for periodontal tissue and investigated the influences of tissue optical properties (anisotropy factor, tissue refraction index and etc.) on energy deposition within the medium. In another case, Parvin *et al.* [19] focused on the energy distribution in dermis using Monte Carlo simulation and studied the effect of anisotropic scattering on the density of absorbed photon. Most previous studies on LLLT considered the spectral range from 600 nm to 1000 nm due to the high penetration depths associated and availability of inexpensive laser diodes within this spectral range.

In recent years, there have been significant advancements in near infrared (NIR) laser technology making inexpensive lasers in the spectral range from 1000 nm to 1900 nm available. In addition, this underutilized wavelength range can provide advantages over the more conventional range, coupling energy more effectively to the tissue for therapeutic applications. However, achieving large enough penetration depths can be challenging and wavelength selection can significantly affect the location of treatment and the dosage of radiative energy deposited. Thus, there is a need for understanding (i) the penetration of these NIR wavelengths in different skin layers as well as (ii) the dependence of the energy deposition within the difference skin layers on control parameters such as wavelength, beam profile and laser power.

Moreover, wavelength selection in LLLT can be optimized using inverse methods for a particular therapeutic application. Use of inverse methods in extracting optical properties is widely reported in the literature [20-24]. For example, Bashkatov *et al.* [22] measured the optical properties of the human skin and mucous tissues by coupling spectrophotometer results to inverse adding-doubling method. They considered human skin as a single-layer medium and reported both absorption and scattering coefficients in the spectral range from 400 nm to 2000 nm. Moreover, Rajaram *et al.* [23] presented a lookup table-based inverse model for determining the optical properties of turbid media from steady-state diffuse reflectance in the wavelength range from 400 nm to 700 nm. More recently, Hennessy *et al.* [24] performed the inverse Monte Carlo method to generate a lookup table for extracting optical properties from tissue-simulating phantom.

Consequently, these inverse methods can be employed to enhance the LLLT results by optimizing the wavelength selection based on the desired penetration depth.

## **1.2 ORGANIZATION OF THE THESIS**

In the first part of the study, Chapter 2, the laser propagation at different wavelengths was compared using reduced-variance Monte Carlo method. In this section skin was modeled as three-layer participating medium with variable optical properties at different layers. Moreover, the skin damage threshold in this wavelength range was investigated and the maximum permissible exposure duration was reported based on ANSI standards.

In Chapter 3, a reduced-variance inverse Monte Carlo method was developed to extract the optical properties of the single-layer skin, based on the intended reflectance and transmittance. The implemented simulation tool was then verified using the available benchmark solutions. Finally, a lookup was generated to provide a scale for selection of wavelength in the spectral range from 1000 nm to 1900 nm.



## **Chapter 2**

# **Light Propagation through the Human Skin; Forward Monte Carlo Method**

This chapter presents a numerical study on laser propagation in human skin using a reduced-variance Monte Carlo method. The wavelength range from 1000 nm to 1900 nm has been considered and the energy deposition within skin layers in this spectral range is explained accordingly.

### **2.1 ANALYSIS**

To analyze the light propagation through a human tissue, a multilayer skin model was considered and laser penetration was traced by Radiative Transport Equation (RTE). Then, the Monte Carlo method was employed to solve the RTE equation for this multilayer medium.

#### **2.1.1 Multi-layer skin model**

In this study, human skin was modeled as a three layer rectangular cuboid with the specific geometry as given in Figure 2.1. The thicknesses of epidermis, dermis and subcutaneous were 0.3 mm, 1.2 mm and 3.0 mm, respectively [25]. A collimated laser beam was incident normally on the top surface of the skin (epidermis). The laser beam diameter was obtained from VetLase Inc<sup>™</sup> and considered to be 6.1 mm.

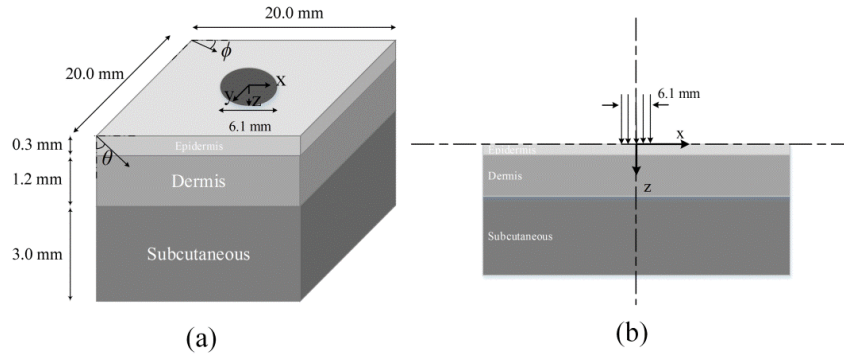


Figure 2.1: The schematic view of three layer skin model:  
 (a) 3-D and (b) 2-D

In order to make the problem mathematically tractable, the following assumptions were made: (1) The refractive index and asymmetry factors were constant within each layer. (2) Surrounding air had the refractive index of 1.00 and was non-participating. (3) Due to the high absorptivity of the muscle layer beneath the subcutaneous [26,27], this layer was considered radiatively cold and black. (4) The blackbody emission by skin and surrounding tissues within the wavelength range of interest was negligible at the normal body temperature of 36.5°C.

### 2.1.2 Governing equations and boundary conditions

Transport of laser radiation within skin can be described by the Radiative Transport Equation (RTE). The three-dimensional steady-state RTE [28] can be written as,

$$\frac{dI_\lambda(S, \Omega)}{dS} = -\kappa_\lambda I_\lambda(S, \Omega) - \sigma_\lambda I_\lambda(S, \Omega) + \frac{\sigma_{s,\lambda}}{4\pi} \int_{4\pi} I_\lambda(S, \Omega_i) \Phi_\lambda(\Omega_i, \Omega) d\Omega_i \quad (2.1)$$

where  $S$  is the path length,  $I_\lambda(S, \Omega)$  is the local radiation intensity,  $\kappa_\lambda$  and  $\sigma_\lambda$  are the absorption and scattering coefficients of skin, respectively. Moreover,  $\Phi_\lambda(\Omega_i, \Omega)$  is the scattering phase function which describes the probability of redirecting the incoming intensity from an arbitrary direction  $\Omega_i$  to the direction of interest  $\Omega$ . In this study, the scattering phase function is approximated with Henyey-Greenstein function [28] as,

$$\Phi_{HG}(\Omega) = \frac{1 - g^2}{(1 + g^2 - 2g \cos(\Omega))^{3/2}} \quad (2.2)$$

where,  $g$  known as the asymmetry factor, is the mean cosine of scattering angle  $\Theta$  and related to the phase function by,

$$g = \overline{\cos(\Theta)} = \frac{1}{4\pi} \int_{4\pi} \Phi(\Theta) \cos(\Theta) d\Omega_i \quad (2.3)$$

Asymmetry factor approaches +1 for strongly forward scattering media, 0 for isotropic scattering, and -1 for strongly backward scattering. In this study the asymmetry factor and the refractive index of each layer has been considered individually and are obtained from literature [29,30], over the spectral range from 1000 nm to 1900 nm. The asymmetry factors considered for epidermis, dermis and subcutaneous were 0.80, 0.91 and 0.75 and the obtained refractive indexes for these layers reported as 1.34, 1.40 and 1.44, respectively.

At the top boundary, reflectance was taken into account through the Fresnel equation according to [31],

$$\rho_{\lambda} = \frac{1}{2} \left[ \left( \frac{n_{\lambda,1} \cos \theta_2 - n_{\lambda,2} \cos \theta_1}{n_{\lambda,1} \cos \theta_2 + n_{\lambda,2} \cos \theta_1} \right)^2 + \left( \frac{n_{\lambda,1} \cos \theta_1 - n_{\lambda,2} \cos \theta_2}{n_{\lambda,1} \cos \theta_1 + n_{\lambda,2} \cos \theta_2} \right)^2 \right] \quad (2.4)$$

where  $\theta_1$  and  $n_{\lambda,1}$  are the angle and the refractive index of the incident side and  $\theta_2$  and  $n_{\lambda,2}$  are referred to those for the transmitted side. Here, the refraction was accounted for by Snell's law [28] as,

$$n_{\lambda,1} \sin \theta_1 = n_{\lambda,2} \sin \theta_2 \quad (2.5)$$

Finally, in the case of collimated radiation, the boundary conditions for the top surface can be written as,

$$I_{\lambda}(x, y, z = 0) = G_{\lambda,0} \delta(\theta_c) \quad (2.6)$$

and for bottom surface,

$$I_{\lambda}(x, y, z = l, \pi/2 < \theta < \pi, \phi) = 0 \quad (2.7)$$

where  $l$  is the thickness of the skin and  $G_{\lambda,0}$  is the incident laser irradiance which can be calculated based on the laser power used.

### 2.1.3 Closure laws

First, the radiation characteristics of the skin layers in the wavelength range of 1000 nm to 1900 nm were compiled from the results reported by Simpson *et al.* [27], Troy and Thennadil [29] and Iino *et al.* [32]. Figure 2.2 demonstrates the absorption and scattering coefficients of the skin layers in this wavelength range.

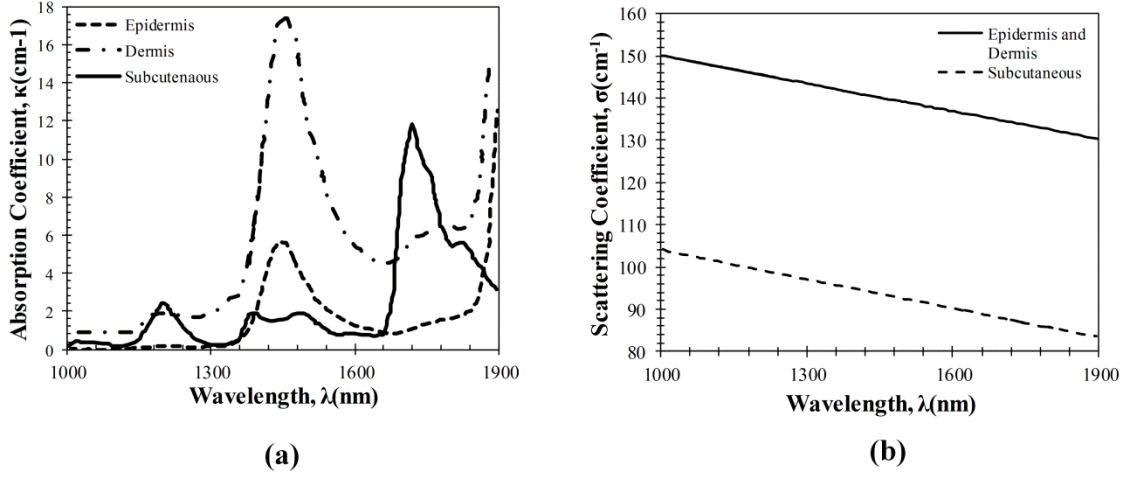


Figure 2.2: (a) Absorption and (b) scattering coefficients of human skin layers in the wavelength range from 1000 nm to 1900 nm

Then, a non-dimensional parameter, namely transport optical thickness, was used to select wavelengths of interest. Transport optical thickness indicates the overall strength of the medium in attenuating radiation at a given wavelength in the penetration direction. It can be written as,

$$\tau_{\lambda} = \sum_{i=1}^n d_i \kappa_i + \sum_{i=1}^n (1 - g_i) d_i \sigma_i \quad (2.8)$$

where  $d_i$  is the physical thickness,  $n$  is number of layers,  $g$  is the asymmetry factor,  $\kappa_{\lambda}$  and  $\sigma_{\lambda}$  are absorption and scattering coefficients of the  $i^{th}$  layer, respectively. Figure 2.3 illustrates this parameter for the wavelengths between 1000 nm to 1900 nm. As inadequate penetration is an issue at these larger wavelengths, wavelengths that corresponded to relatively low transport optical thicknesses have been considered, as indicated.

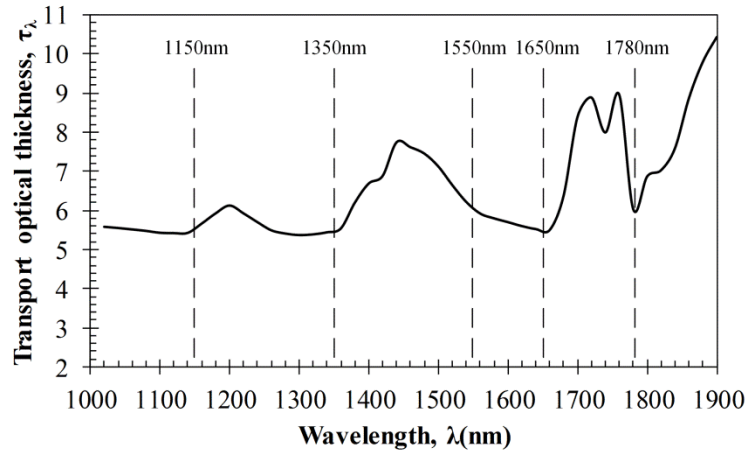


Figure 2.3: Transport optical thickness of human skin for wavelengths between 1000 nm to 1900 nm

### 2.1.3 Method of solution

In this study, a three-dimensional multi-layer reduced-variance Monte Carlo method was implemented and validated against benchmark solutions. The Monte Carlo technique was first proposed by Metropolis and Ulam [33] as a statistical approach for studying differential equations. More recently, Prahl *et al.* [16] applied the variance reduction technique, proposed by Kahn and Harris [34], to the Monte Carlo method and implemented a reduced-variance Monte Carlo for simulating light propagation in tissue. By using the variable step-size walking process in this method, the high variance in the normal Monte Carlo method can be resolved without a considerable CPU time increase. To illustrate, light is modeled as a large number of packets that carry a specific amount of energy (i.e., 1.0). The method tracks each energy packet as it is absorbed, scattered, back-scattered, transmitted or reflected within the skin. For the packets which scatter within the skin, a fraction of their energy will be absorbed and deducted at each interaction as,

$$W = \frac{\kappa_\lambda}{\kappa_\lambda + \sigma_\lambda} W_0 \quad (2.9)$$

where  $W_0$  is the packet's initial energy, and  $W$  is the packet's energy after the interaction. Based on this method, a packet can propagate infinitely and its energy will never reach zero. Thus, an energy threshold is defined to terminate the packet with an energy level below the threshold. This threshold energy level was optimized to ensure that the results are independent of this selection so energy of 0.001 was considered for this study. In order to track late-stage propagation, such a weakened packet is given a second chance in  $m^{\text{th}}$  (i.e.,  $m=10$ ) stage of surviving through roulette technique. According to this technique, if a packet was not absorbed by the ninth interaction with the medium, the roulette will either empower it to  $mW$  or eliminate it, on a random basis. The killed packets are monitored to satisfy the energy conservation. Then, based on the statistical results of the energy packets, the absorptivity and transmittance for the considered skin can be calculated. Figure 2.4 shows the flow chart for this multi-layer three-dimensional reduced-variance Monte Carlo scheme used in this study.

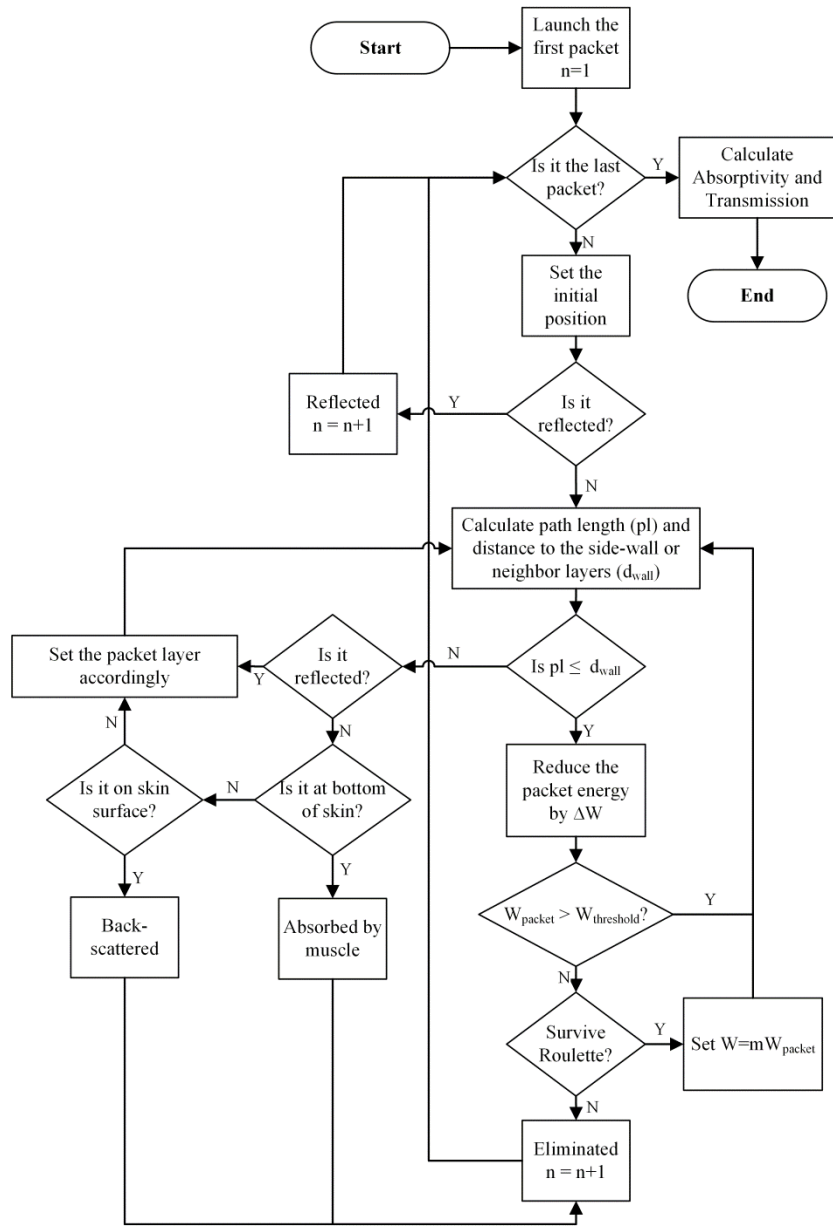


Figure 2.4: The flow chart of solution algorithm used in this study for the multi-layer three-dimensional reduced-variance Monte Carlo scheme

This code was validated for both scattering and non-scattering media cases. In the non-scattering case, a rectangular cuboid of water has been considered as a purely absorbing medium. Transmissivity of the cuboid layer was obtained from Monte Carlo



method as well as Beer-Lambert's law [28], using Wieliczka *et al.* [35] data for water optical properties. In the case of anisotropic media, a medium with the optical thickness of 2.0, albedo of 0.9 and asymmetry factor of 0.75 was simulated and the obtained results were compared to those of Van de Hulst [36] and Prahl *et al.*[16]. Based on these comparisons, the percent deviations for the transmittance with respect to the Beer-Lambert law, Van de Hulst's tables and Prahl *et al.* were 0.01%, 0.11% and 0.30%, respectively.

Moreover, the sensitivity of the Monte Carlo code on the number of energy bundles used was analyzed to ensure that the obtained results were independent of this choice. By increasing the number of packets from  $10^6$  to  $10^7$ , changes in the absorptance results were less than 0.1%, but the CPU time increased by 20 hours. Thus,  $10^6$  energy bundles with the total power of 155mW were used in all of the results reported in this chapter.

## **2.2 RESULTS AND DISCUSSION**

First the reduced-variance Monte Carlo simulation at the wavelength of the 1350 nm was operated to investigate the light propagation phenomena. In the next step, the simulation has been performed at the selected wavelengths to study the penetration and absorption of lasers with uniform profiles in different skin layers. Then, the results were processed in the perspective of power delivery and the selected wavelengths were compared based on this factor. Moreover, simulations for different Gaussian beam profiles were performed and dependence of laser penetration and absorption in skin on

the beam profile characteristics has been shown. Lastly, the maximum allowable exposure time has been reported for the considered laser probe based on the maximum absorbed fluence and available standards.

### 2.2.1 Light propagation analysis at the wavelength of $\lambda = 1350$ nm

The relevant parameter in laser treatment of tissues is the dosage, i.e., the amount of energy density absorbed by tissue. Local variations of dosage can significantly affect a treatment and can be defined as,

$$D_v = P_{abs} \cdot t_{exp} \quad (2.10)$$

where  $P_{abs}$  is the volumetric absorbed power density ( $mW/mm^3$ ) and  $t_{exp}$  is the exposure duration. Thus, here the results have been reported based on the density of absorbed power which provides the flexibility of calculating the dosage at different exposure times.

For illustrative purposes, a collimated laser beam of diameter 6.1 mm has been considered with irradiance of 155 mW at the wavelength of 1350 nm. The local absorbed power density was measured by using the results of the implemented Monte Carlo code. To do so, for a given volume, the fraction of absorbed packets were counted and multiplied to the amount of energy that each packet carries. Hence, by dividing the measured absorbed energy by its corresponding volume, the local absorbed power density was measured. Figure 2.5 shows the volumetric power density as a function of location within the three layer skin model on the  $x-z$  plane. In order to generate this map, the 3D

results of absorbed power within the skin were integrated in  $y$  - direction and projected on to the  $x-z$  plane. Besides, in this figure, the percentages of the absorbed power at each layer are also shown. These percentages show the absorptance of laser within the skin layers so the rest of the energy packets can be either transmitted or back-scattered from the skin surface.

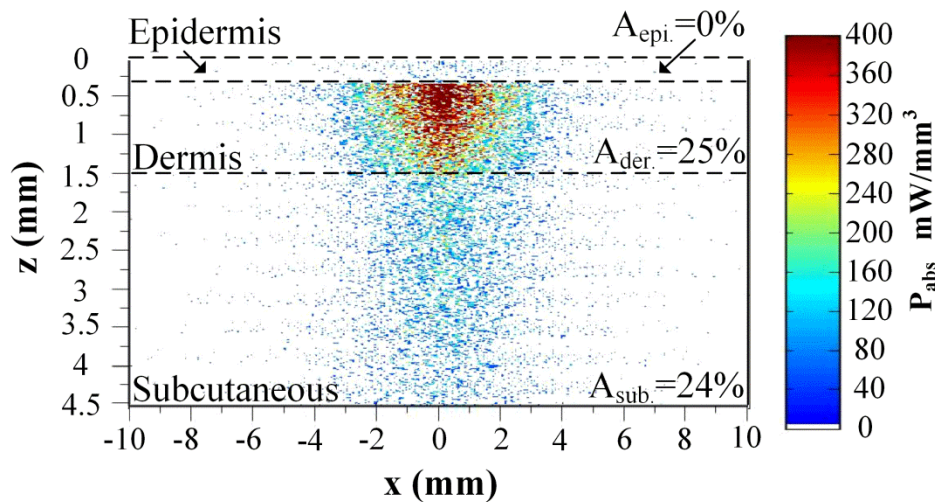


Figure 2.5: Location and volumetric density of absorbed power for  $P_{in}=155$  mW at wavelengths of 1350 nm. Absorptance of each layer is stated in the figure in percentages;  $x$  describes the location on  $x$ ;  $z$  represents the distance from the surface and colors indicate the volumetric density of absorbed power within the skin

As it can be seen in Figure 2.5, the absorbed power within the epidermis layer is negligible compared to those within dermis and subcutaneous. Moreover, although the absorptance of dermis and subcutaneous are fairly close, the propagation of the absorbed power in these two layers is different. In dermis, the dosage of absorbed power reached the highest level of  $400 \text{ mW/mm}^3$ , whereas in subcutaneous it never exceeded  $200 \text{ mW/mm}^3$ . To see the contribution of each layer to the absorption more accurately, Figure

2.6 shows the local absorbed power of skin layers as a function of distance from the skin surface. This figure shows that the local absorptance decreases for dermis and subcutaneous as the light penetrates. Based on this figure, 56% and 48% decrease in local absorptance was observed within the layers of dermis and subcutaneous, respectively. Moreover, Figure 2.6 confirms that the level of absorptance is considerably higher in dermis compared to other layers, as it was also showed in Figure 2.5.

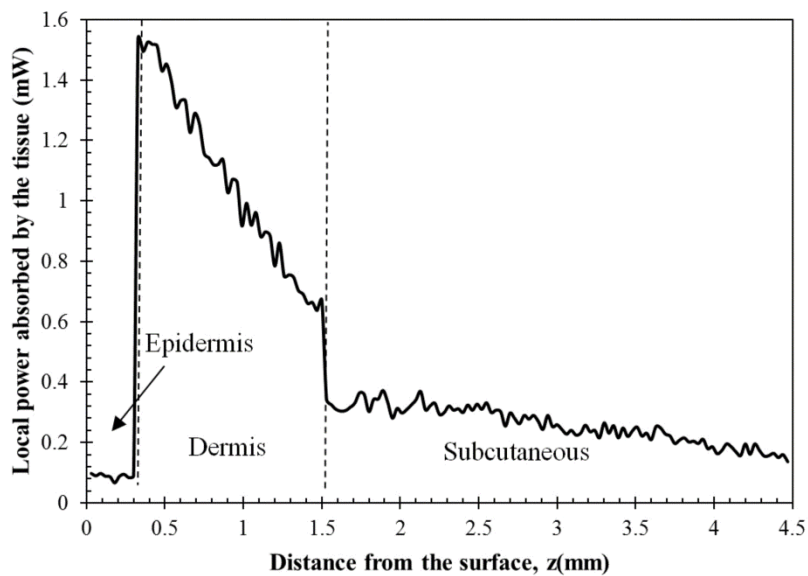


Figure 2.6: Local power absorbed within the skin as a function of distance at  $\lambda=1350$  nm and  $P_{in} = 155$  mW

### 2.2.2 Penetration at selected wavelengths

As it was discussed previously, in order to have the most effective result in LLLT, the laser power should be efficiently delivered to the skin layer of interest. Based on the purpose of therapy, this layer can be epidermis, dermis, subcutaneous or the muscle layer beneath the subcutaneous.

In this section, the penetration of the light at the selected wavelengths (Section 2.1.3) has been compared based on their local absorptance within their layers. To illustrate, the location of the absorbed power density within the skin has been plotted for these wavelengths in Figure 2.7. These results indicate that although the total absorptance of the skin is similar at 1150 nm, 1350 nm, and 1650 nm, the spatial distribution of absorbed power within the different layers of skin is quite different. Particularly at 1350 nm, the incident laser power is mostly delivered to the dermis, 25%, whereas in the cases of 1150 nm and 1650 nm, subcutaneous absorbed the highest percentage of the incident laser power, about 35%. Moreover, the color plots show that the distribution of absorbed power within the skin layers was highly dependent on the wavelength. For example, the volumetric power density was fairly uniform across the dermis and subcutaneous at 1150 nm whereas high volumetric absorption densities were observed in the dermis at wavelengths 1350 nm, 1550 nm, and 1650 nm. Furthermore, Figure 2.7(e) shows that unlike at any other wavelength considered, power at 1780 nm could penetrate through epidermis and dermis and get highly absorbed by subcutaneous layer. At this wavelength 90% of laser power was absorbed by subcutaneous without significant absorption in epidermis and dermis.

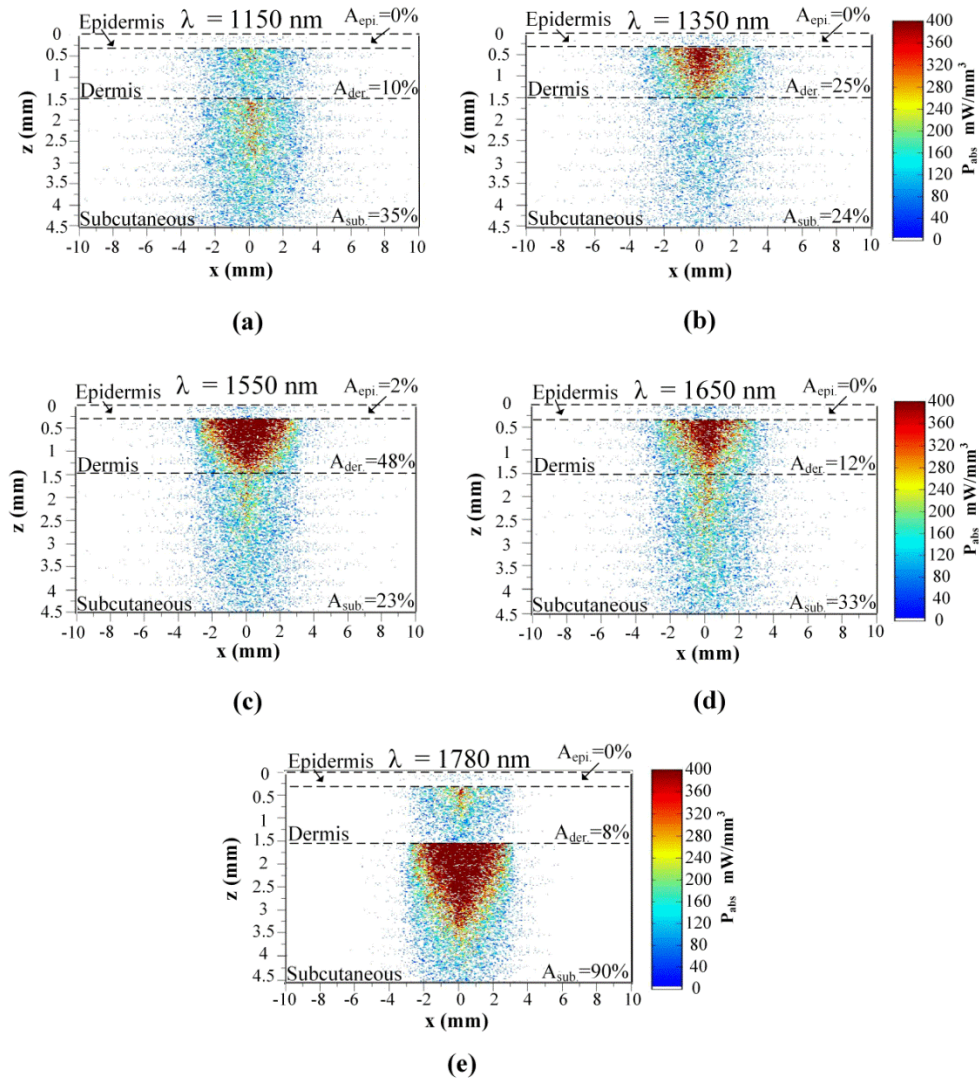
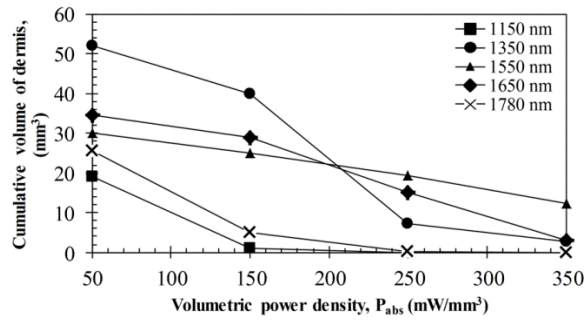


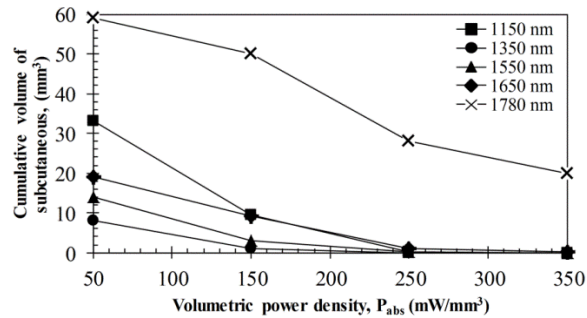
Figure 2.7: Location and volumetric density of absorbed power for  $P_{\text{in}}=155 \text{ mW}$  at wavelengths of (a) 1150 nm, (b) 1350 nm, (c) 1550 nm, (d) 1650 nm and (e) 1789 nm. Absorptance of each layer is stated in the figure in percentages;  $x$  describes the location on  $x$ ;  $z$  represents the distance from the surface and colors indicate the volumetric density of absorbed power within the skin

### 2.2.3 Dosage delivery at selected wavelengths

From the medical application perspectives, it is important to know the volume of tissue absorbing a specific dosage of energy in a given skin layer. In order to bring this point in perspective, Figure 2.8 shows the cumulative volume of skin layers, namely dermis and subcutaneous, that absorbed the power density equal to or larger than the indicated amount. To illustrate, Figure 2.8(a) demonstrates that at 1350 nm, the volume of dermis that was subjected to a power density of  $250 \text{ mW/mm}^3$  and lower, was largest despite the fact that dermis had the largest percentage of incident power absorbed at 1550 nm. However, when the total magnitude of the highest absorptance is concerned, 1550 nm was more prevalent in delivering this range of power to dermis. In addition, Figure 2.8(a) shows that the volumetric treatment of dermis at similar power density were similar for 1150 nm with 1780 nm and 1550 nm with 1650 nm. On the other hand, Figure 2.8(b) shows that 1780 nm featured the largest power density and volume of tissue treated in subcutaneous layer. In this layer, the local power density was lower than  $250 \text{ mW/mm}^3$  for all other wavelengths considered where 1150 nm featured the largest volume of absorption.



(a)



(b)

Figure 2.8: The cumulative volume of (a) dermis and (b) subcutaneous that has the power density equal to or larger than the indicated amount in the horizontal axis

In deep tissue treatments, the total power transmitted through the skin is of interest rather than the local absorptance within the different skin layers. Figure 2.9 shows the total transmitted power as a function of transport optical thickness based on the simulation results in the considered spectral range. This figure indicates that the transmitted power decreases exponentially with the transport optical thickness. Having similar transport optical thicknesses indicate that the same percentage of incident power will be transmitted through the entire skin which can be used as a first order tool for identifying and selecting wavelengths and power for deep tissue LLLT applications. As an example, at wavelength of 1650 nm, which has the transport optical thickness of 5.5,



more than 41% of the laser power can be delivered to deep tissue whereas at 1250 nm, 1580 nm and 1700 nm, with the optical thicknesses of 5.6, 5.8 and 8.4, this transmitted power decreases to 37%, 29% and 1%, respectively.

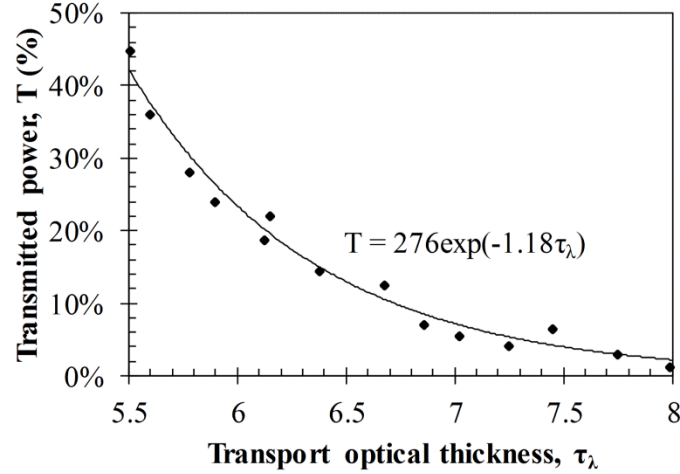


Figure 2.9: Correlation between transport optical thickness and the percentage of power than was transmitted through the entire skin without being absorbed

#### 2.2.4 Effect of beam profile on the light propagation

For many medical laser applications, beam profile is a critical parameter. It can affect the volumetric dosage as well as the penetration depth of the incident laser power. In the previous analysis, a uniform beam profile was investigated, but a Gaussian profile is often encountered in medical applications [37]. For a laser with a Gaussian profile, the irradiance distribution is given as,

$$G_{\lambda,xy} = G_{\lambda,max} \exp\left(-\frac{x^2 + y^2}{\sigma^2}\right) \quad (2.11)$$

where  $G_{\lambda,xy}$  is the local irradiance and  $G_{\lambda,max}$  is the irradiance at the beam center ( $x = 0, y = 0$ ).  $\sigma$  is the divergence radius which indicates the specific annular distance from the center where  $1/e$  of maximum irradiance has been reached. By varying this parameter the laser beam profile can be changed from a very sharp peak when  $\sigma \rightarrow 0$ , to the uniform profile at  $\sigma \rightarrow \infty$ . In order to see the influence of different spatial beam profiles on penetration and absorption, three different Gaussian beam profiles have been considered. The total power of the beams considered were the same in all cases but each feature a divergence radius of 1 mm, 2 mm and 10 mm, respectively. Figure 2.10 shows the irradiance profiles of these Gaussian beams.

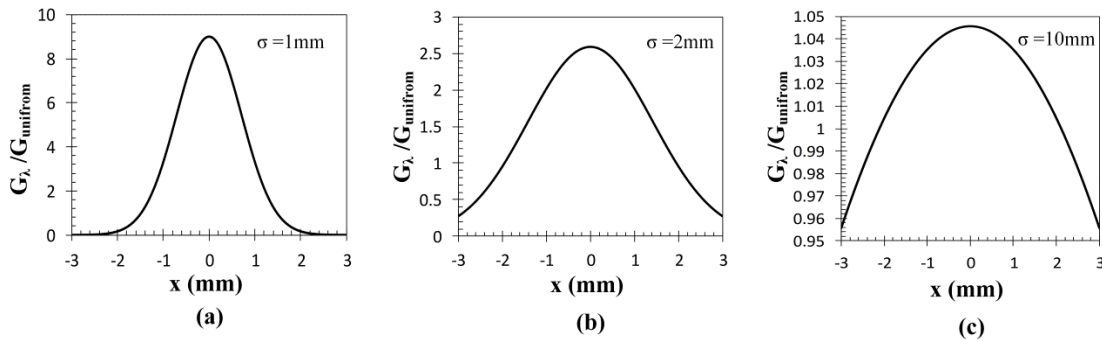


Figure 2.10: Irradiance profile of three Gaussian laser beams with a divergence radius of (a) 1mm, (b) 2mm and (c) 10mm;  $x$  is the distance from the beam center and  $G_{uniform}$  is the irradiance of the uniform profile which carries the same amount of total power.

The propagation of these beams at the wavelength of 1350 nm in skin has been simulated and the results with those of the uniform profile case have been compared. Although the total absorbed and transmitted power through the entire skin was similar, there was a significant difference in the local volumetric dosages. Figure 2.11 shows the

volumetric dosage for different beam profiles at the wavelength of 1350 nm. According to this figure, for the Gaussian beam with divergence radii of 10 mm, the distribution of absorbed power was similar to that of a uniform beam profile. However, smaller divergence radii results in highly localized larger volumetric dosages along the centerline of the beam.

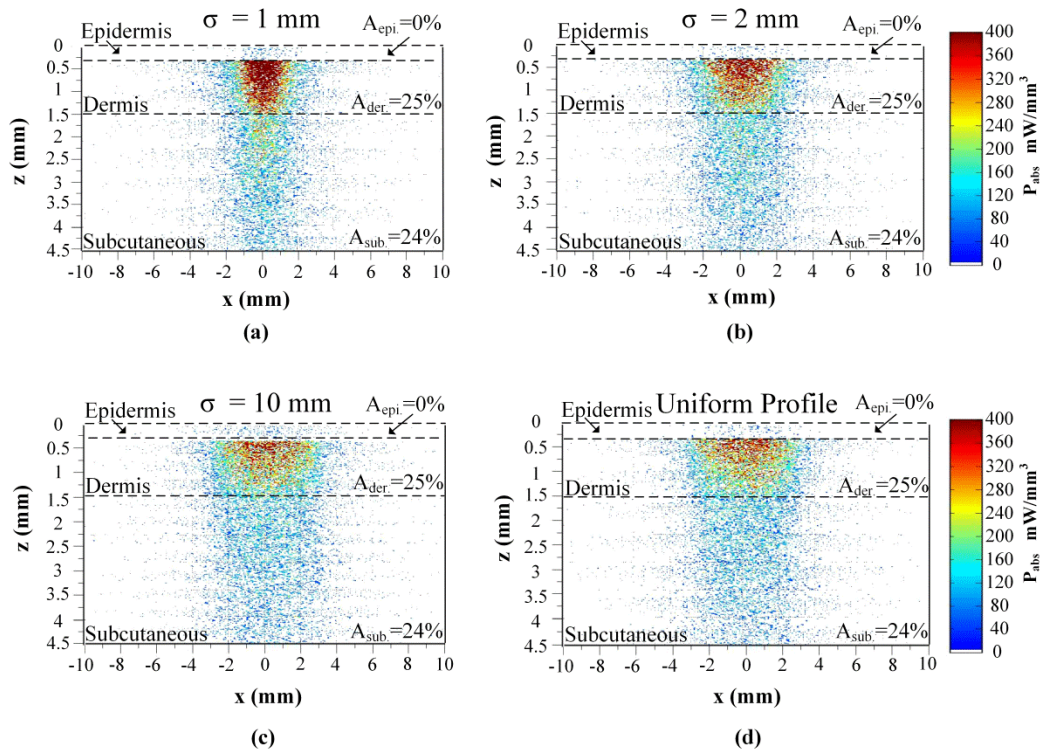


Figure 2.11: Location and volumetric dosage of absorbed energy for  $P_{in}=155 \text{ mW}$ , at the wavelength of 1350 nm, for Gaussian beam profiles with divergence radii of (a) 1mm, (b) 2mm, (c) 10mm, and (d) uniform beam profile

To further illustrate this point, Figure 2.12 shows the volumetric dosage at the two different locations with respect to the beam center, namely at the beam center  $x = 0 \text{ mm}$  and at 2 mm away from the beam center,  $x = \pm 2 \text{ mm}$ . Figure 2.12(a) shows that by

changing the beam profile from uniform to Gaussian, the dosage of absorbed power can be increased from  $150 \text{ mW/mm}^3$  to  $450 \text{ mW/mm}^3$  in dermis along the beam center. Similarly, in subcutaneous, the local dosage was about three times higher for the sharpest Gaussian profile in comparison to uniform beam. However, according to Figure 2.12(b), at  $x = \pm 2 \text{ mm}$ , the local dosage in dermis and subcutaneous decreased from  $120 \text{ mW/mm}^3$  to  $40 \text{ mW/mm}^3$  and from  $30 \text{ mW/mm}^3$  to  $10 \text{ mW/mm}^3$ , respectively.

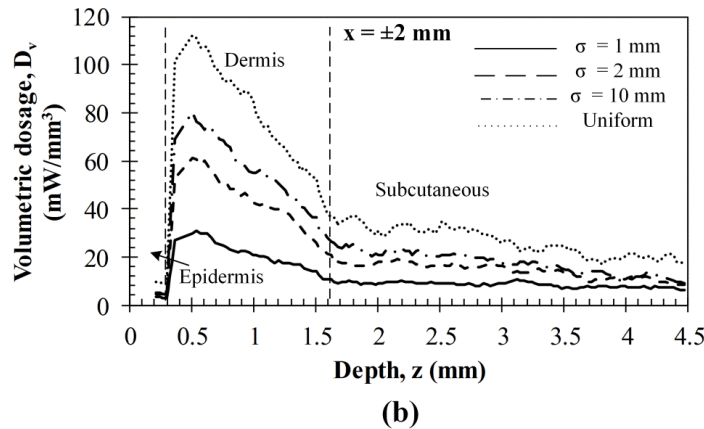
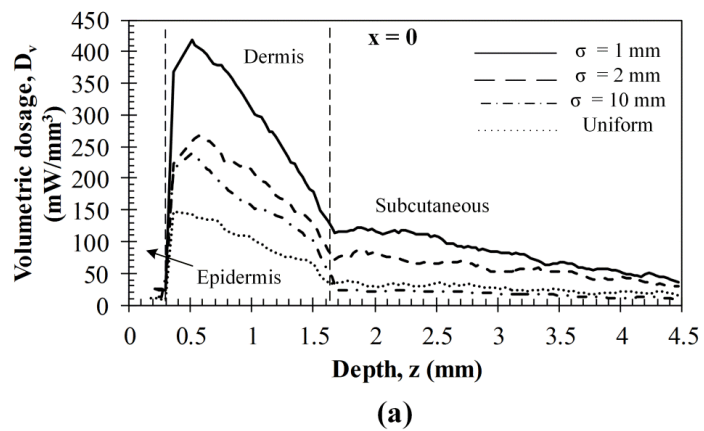


Figure 2.12: The volumetric dosage versus depth from the surface for  $P_{in}=155 \text{ mW}$ , at two locations of (a)  $x=0$  and (b)  $x=2 \text{ mm}$ ;  $x$  is the distance from the center of the beam

### 2.2.5 Exposure duration and skin damage threshold

Skin damage threshold can be defined as a critical amount of absorbed radiation energy, which causes permanent denaturation of skin cells. Above a critical fluence, laser irradiation causes the necrosis of tissue cells and lethal damage to superficial blood vessels [38,39]. In order to prevent this damage, the American National Standards Institute (ANSI) provides the Maximum Permissible Exposure (MPE) per area as a function of duration of exposure for different wavelengths based on experimental results [40]. These MPE data includes visible and near infrared part of the electromagnetic spectrum for wavelengths from 400 nm to 2600 nm. To find the maximum permissible exposure duration of the studied wavelengths, the maximum absorbed fluence within the skin (i.e., absorbed power per area) from the simulations was obtained and the ANSI specifications at the corresponding wavelengths were applied. Table 2.1 shows the maximum permissible exposure duration at the considered wavelengths with the laser power of  $p_{laser}$  expressed in mW. In this table the beam diameter was considered to be 6.1 mm, however, a sensitivity analysis was performed to ensure that the beam diameter does not affect the MPE values by more than 2% in the diameter range from 5 to 15 mm.

Table 2.1: Maximum permissible exposure (MPE) duration for the wavelengths of 1150, 1350, 1550, 1650 and 1780 nm at the power of  $p_{laser}$  (mW) according to ANSI Z136.1-2007 [40]

Wavelength (nm)	Permissible power range (mW)	Maximum permissible exposure duration (s)
1150	$p_{laser} \leq 10^3$	$2.9/p_{laser}$
1350	$p_{laser} \leq 10^3$	$2.3/p_{laser}$
1550	$p_{laser} \leq 700$	$10/p_{laser}$
1650	$p_{laser} \leq 700$	$10/p_{laser}$
1780	$p_{laser} \leq 700$	$10/p_{laser}$

The maximum exposure durations can easily be calculated using Table 2.1 for lasers with typical powers within the ranges of LLLT. For example, for the studied laser (i.e., 155 mW), the maximum exposure durations for the wavelength of 1150 nm and 1350 nm are about 19 s and 15 s, respectively. In addition, for higher wavelengths of 1550 nm, 1650 nm and 1780 nm, this duration can be increased up to 64 s. Although no damage in tissue is guaranteed, it should be noted that these exposure durations are on the conservative side.

### 2.3 CONCLUSION

In this chapter, in order to study the laser penetration through human skin, this tissue was modeled as a three-layer participating medium. A collimated laser was then incident on the top layer and penetrated through the skin. Radiative Transport Equation (RTE) was employed to simulate the laser propagation within the skin. In order to select the wavelengths of interest a dimensionless number, namely transport optical thickness, was used. Wavelengths with the lowest optical thickness were selected to maximize the

penetration depth and absorbed dosage. To solve the RTE for these wavelengths, a three-dimensional multi-layer reduced-variance Monte Carlo method was implemented and validated against the benchmark solutions. This Monte Carlo method was used to evaluate the effect of wavelength and beam profile on the laser penetration through the skin. In addition, ANSI standards were employed to calculate the maximum permissible exposure duration for the considered wavelengths and laser power. Based on the results obtained, the following conclusions can be drawn:

- Based on the selected wavelength, the laser power can be efficiently delivered to the specific layer of the skin. For example, 90% absorptance in subcutaneous and 48% absorptance in dermis can be obtained by using wavelengths of 1780 nm and 1550 nm, respectively.
- Although the total absorptance at the wavelengths of 1150 nm, 1350 nm and 1650 nm are fairly close, the spatial distribution of the absorbed power within the skin layers is considerably different. To illustrate, dermis is the main absorber at the wavelength of 1350 nm, whereas in the cases of 1150 nm and 1650 nm, the incident power was mostly deposited to subcutaneous.
- Similar absorptance rate in one layer does not indicate the similar laser propagation. For example, absorptance level in dermis is fairly similar at both wavelengths of 1150 nm and 1650 nm. However, the volumetric dosage in dermis at wavelength of 1150 nm never exceeded the value of  $280 \text{ mW/mm}^3$ , while at wavelength of 1650 nm this value was as high as  $400 \text{ mW/mm}^3$ .

- In case of power delivery to dermis, wavelength of 1350 nm for dosages lower than  $150 \text{ mW/mm}^3$  and 1550 nm for dosages higher than  $250 \text{ mW/mm}^3$  should be utilized. In addition, wavelength of 1780 nm featured the highest power delivery to subcutaneous.
- At a given laser power, the local dosage of absorbed power can be increased about three times, by changing the beam profile from uniform (flat) to a sharp Gaussian's one.
- Maximum permissible exposure durations is inversely related to the laser power. For the laser power of 155 mW, the exposure durations can be up to 19 s and 15 s at wavelengths of 1150 nm and 1350 nm, respectively.



## Chapter 3

### Penetration depth correlation with optical properties; the Inverse Monte Carlo method

As discussed in Chapter 2, based on the selected wavelength, the laser power can be either absorbed within the skin layers or in muscle layer beneath the skin. For different purposes of the therapy, a particular combination of these two absorptions might be needed. For example, for deep tissue treatments the laser power is preferably deposited to the muscle layer while in hair growth, dermis is the target layer of the therapy. These two absorptions can be quantified by measuring Transmittance ( $T_\lambda$ ), Reflectance ( $R_\lambda$ ) and Absorption ( $A_\lambda$ ) for a given therapy. Transmittance indicates the portion of the energy that is deposited to the muscle layer beneath the skin, absorptance stands for the portion that absorbed within the skin and reflection represents the fraction of the reflected energy. Considering that there is a negligible energy loss from sides of the medium, the following equation can be derived,

$$T_\lambda + R_\lambda + A_\lambda = 1 \quad (3.1)$$

Thus, by knowing two of these energy fractions, a therapy can be completely specified. Due to complexity of measuring the absorptance experimentally, transmittance and reflectance are mostly measured and reported through the experiments. Consequently,

this chapter is focused on evaluating the light penetration based on transmittance and reflectance. First, an inverse Monte Carlo method is implemented to recover the radiation characteristics (i.e. absorption and scattering coefficient) from experimentally measured spectral transmittance and reflectance. Then, the correlation between transport optical thickness and values of transmittance and reflectance is established, using the recovered absorption and scattering coefficients. Finally, in order to investigate the effect of wavelength on penetration depth, this depth was measured and plotted in the wavelength range from 1000 nm to 1900 nm, using the forward Monte Carlo for the recovered optical properties.

### **3.1 ANALYSIS**

To recover the absorption and scattering coefficient, a single-layer skin model was considered in this chapter. Then, an inverse Monte Carlo method was generated to retrieve these radiation characteristics from experimental values of transmittance and reflectance.

#### **3.1.1 Single-layer skin model**

Figure 3.1 shows the schematic of the skin model used in this chapter. The thickness of the skin is considered to be 4.5 mm and the laser characteristics are the same as described in Section 2.1. The collimated laser penetrates through the z-direction and propagates laterally in x-y plane.

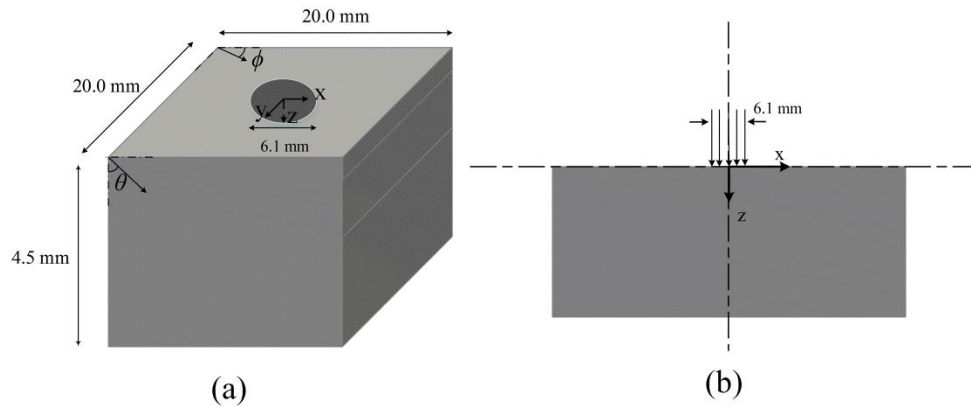


Figure 3.1: The schematic view of single-layer skin model:  
 (a) 3-D and (b) 2-D

In order to make the inverse problem mathematically traceable, the following assumptions were considered: (1) The refractive index and asymmetric factor are assumed to be constant within the considered wavelength range with values of 1.37 and 0.9, respectively. (2) Surrounding air had the refractive index of 1.00 and was non-participating. (3) Due to the high absorptivity of the muscle layer beneath the subcutaneous, this layer was considered radiatively cold and black [26,27]. (4) The blackbody emission by skin and surrounding tissues within the wavelength range of interest (i.e. 1000 nm to 1900 nm) was negligible at the normal body temperature of 36.5°C.

### 3.1.2 Method of solution

A single-layer reduced-variance inverse Monte Carlo method was implemented to find the required wavelength for obtaining the desired reflectance ( $R_\lambda$ ) and transmittance ( $T_\lambda$ ). In this method, for a given reflectance and transmittance, an initial set of optical

properties (i.e. absorption and scattering coefficients) were assumed as the first guesses. Then, using these values, the forward Monte Carlo code (discussed in section 2.1.3) was performed to determine the reflectance and the transmittance. Next, the sum of the squared errors ( $e_{\text{model}}$ ) between the predicted and intended values was calculated using,

$$e_{\text{model}} = \left[ \left( \frac{R_{\text{model}} - R_{\lambda}}{R_{\lambda}} \right)^2 + \left( \frac{T_{\text{model}} - T_{\lambda}}{T_{\lambda}} \right)^2 \right]^{1/2} \quad (3.2)$$

where,  $R_{\text{model}}$  and  $T_{\text{model}}$  are the predicted values and  $R_{\lambda}$  and  $T_{\lambda}$  are the target values for reflectance and transmittance, respectively. The scattering and absorption coefficients are then iteratively updated to minimize the error using the interior point optimization method [41]. To employ this method, a reflectance-transmittance lookup table was generated. This table correlates the values of the reflectance and transmittance to their corresponding absorption and scattering coefficient within the wavelength range from 1000 nm to 1900 nm with the increments of 5 nm. Thus, a total of 180 forward Monte Carlo simulations were needed to create this lookup table. At each iteration, this lookup table was used to find the closest wavelengths ranges. Then, cubic splines were utilized to interpolate between the lookup table values and find the exact wavelength. This loop was repeated until the error,  $e_{\text{model}}$ , was lower than 5%. Figure 3.2 demonstrates the flowchart of the algorithm that used in this method.

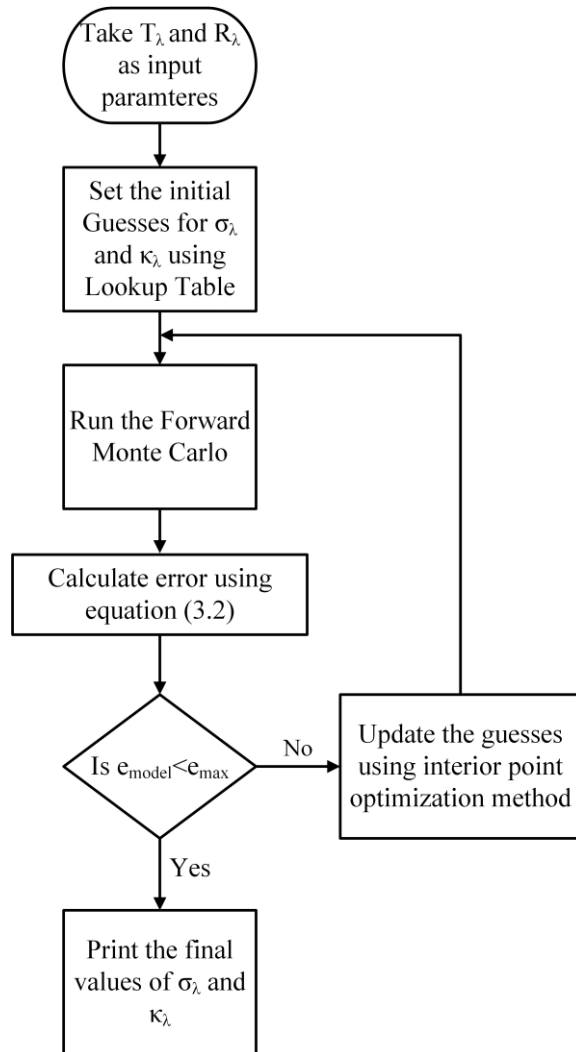


Figure 3.2: The flow chart of solution algorithm used in this study for finding the required wavelength based on given reflectance and transmittance

The forward Monte Carlo method used in this inverse problem is the single-layer version of the reduced-variance multi-layer three-dimensional Monte Carlo, described in section 2.1.3.

## **3.2 RESULTS AND DISCUSSION**

In this section, the accuracy and validation of the implemented inverse method is discussed first. Then, by using this verified inverse Monte Carlo tool, optical properties of the skin is recovered from the experimental values of transmittance and reflectance, obtained from literature. Finally, the correlation between transport optical thickness and experimental values of transmittance and reflectance is studied and the penetration depth of laser in LLLT is discussed as well.

### **3.2.1 Validation and Accuracy Analysis**

To evaluate the performance of the implemented inverse Monte Carlo model, the simulation results were validated against the benchmark solutions. For a medium with the asymmetry factor of 0.75, transmittance of 66% and reflection of 0.97%, the obtained optical properties of the implemented method was verified, using the values that reported by Prahl *et al.*[17] and van de Hulst [36]. Compared to these references, the percent deviation for the absorption and scattering coefficients were measured to be 1.8% and 1.1%, respectively. In the case of isotropic scattering, the simulations results were compared to those of Giovanelli [42] and Prahl *et al.* [17] for a semi-infinite slab with total reflection of 26% and index refraction of 1.5. For this case, the percent deviation of the optical properties was lower than 2% for both scattering and absorption coefficient. Moreover, to ensure that the assumptions of constant asymmetry factor and refractive index do not affect the results considerably, a sensitivity analysis was performed. Based on this analysis, variation of the refractive index within the range of 10%, changes the

values of transmittance and reflectance for no more than 1.5%. Also, a variation of 0.5% was observed by changing the asymmetry factor within the range from 0.85 to 0.95.

Furthermore, to test the accuracy of the implemented inverse Monte Carlo method, the code was utilized to determine the optical properties of the mucous membrane of the maxillary sinuses in the wavelength range from 400 nm to 2000 nm. For this simulation, the mucous membrane was modeled as a semi-infinite slab with the area of  $20 \times 20 \text{ mm}^2$ , asymmetry factor of 0.9 and refractive index of 1.45. Figure 3.3 compares the results of the inverse Monte Carlo method and the experimental results of Bashkatov *et al.* [22] for both absorption and scattering coefficients. In this figure, to better illustrate the differences between simulation results and experimental ones, reduced scattering coefficient was used instead of scattering coefficient which can be defined as,

$$\mu'_s = \sigma_\lambda(1-g) \quad (3.3)$$

where,  $g$  is the asymmetry factor and  $\sigma_\lambda$  is the scattering coefficient of the medium. As it can be seen in this figure, the implemented method provides an excellent fit with the experimental results in both scattering and absorption coefficients. Based on these results, the inverse Monte Carlo method predicted the optical properties over the range from 400 nm to 2000 nm with the root-mean-square percent errors of 0.63% for absorption coefficient and 1.0% for reduced scattering coefficient.

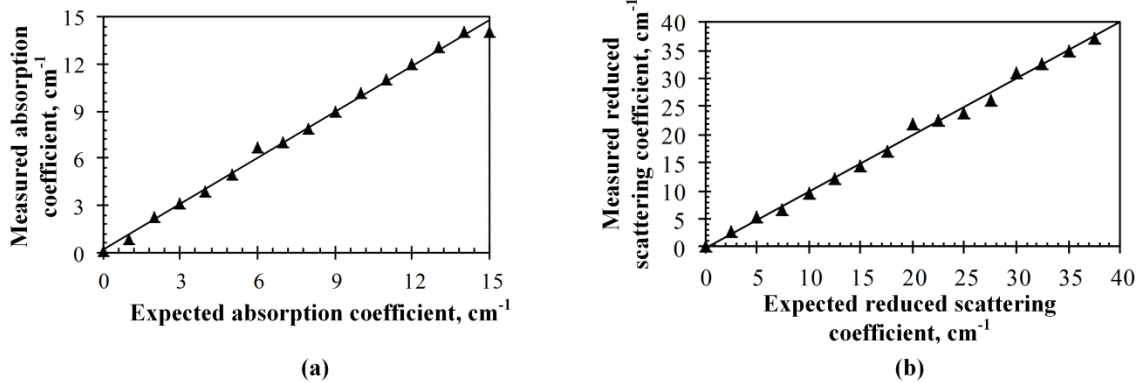


Figure 3.3: Benchmark values [22] versus simulation results for (a) absorption coefficient (b) reduced scattering coefficient of human skin in wavelength range from 400 nm to 2000 nm. The solid line indicates the perfect agreement.

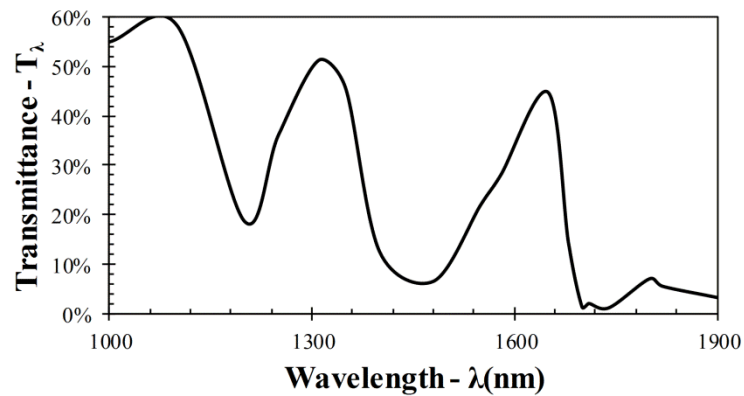
Finally, since the reduced-variance Monte Carlo method was used in this inverse method, the variance of the obtained results was lower than 0.6% for all the cases. By increasing the number of energy packets from  $10^6$  to  $10^7$  the variance of the results can be decreased to about 0.1%, but the CPU time increased about 10 times for each iteration. Thus,  $10^6$  packets were considered for all the results reported in this chapter.

### 3.2.2 Retrieve of optical properties

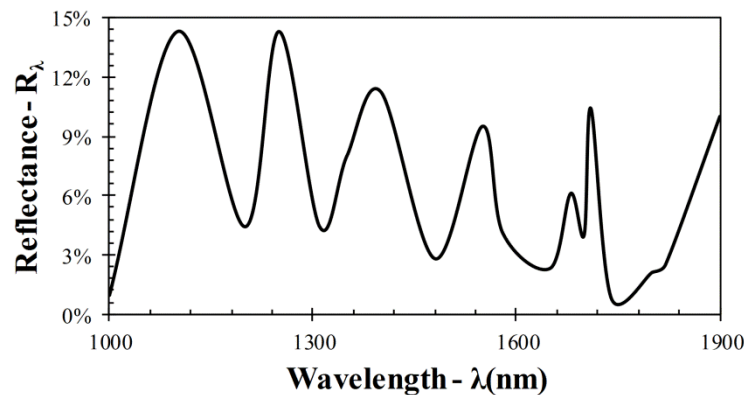
To recover the optical properties of the skin, the experimental results of Hardy *et al.* were employed [43]. They reported the spectral values of reflectance and transmittance, using the absolute method of goniometric spectrophotometry. They obtained the excised skin samples either from surgical specimens or from autopsy and performed the experiments in the wavelength range from 600 nm to 2400 nm. Figure 3.4 shows their results for reflectance and transmittance of the human skin in the spectral



range from 1000 nm to 1900 nm. This figure demonstrates that transmittance and reflectance are highly sensitive to wavelength variations. It is interesting to note that both transmittance and reflectance showed a dynamically increasing and decreasing behavior within the wavelength range from 1000 nm to 1900 nm. Moreover, the periods of these fluctuations were different for transmittance and reflectance.



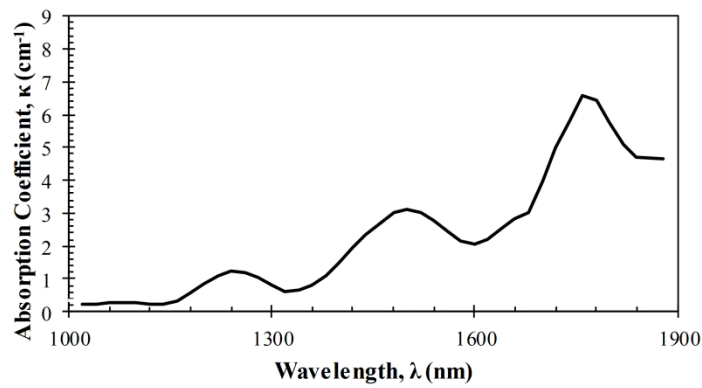
(a)



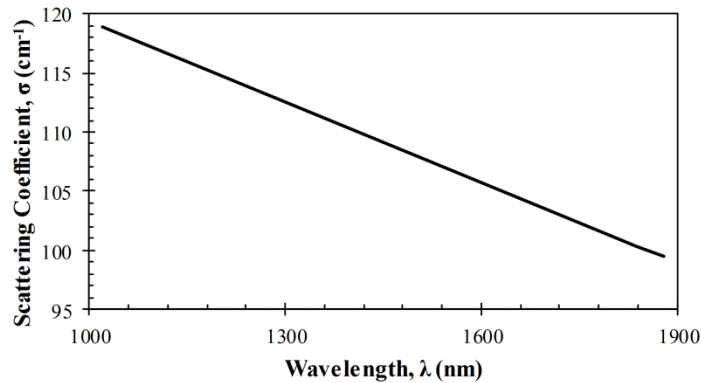
(b)

Figure 3.4: Variations of experimentally obtained (a) transmittance and (b) reflectance of human skin within the wavelength range from 1000 nm to 1900 nm [43]

In the next step, by using the experimental data given in Figure 3.4, the optical properties of the skin were obtained. Figure 3.5 shows the absorption and scattering coefficient of the human skin in this considered spectral range. To generate this figure, 60 inverse Monte Carlo simulations were performed within the range of 1000 nm to 1900 nm with the increments of 15 nm. Intermediate values were interpolated from these simulation results using the third-order polynomial interpolation.



(a)



(b)

Figure 3.5: (a) Absorption and (b) scattering coefficients of human skin recovered from experimental values of transmittance and reflectance, using inverse Monte Carlo method

Then, by using the obtained absorption and scattering coefficients, the transport optical thickness can be simply measured based on Equation 2.8. To recall, transport optical thickness indicates the overall strength of the medium in attenuating radiation at a given wavelength in the penetration direction. Figure 3.6 illustrates the transmittance and reflectance as a function of transport optical thickness. According to this figure, both reflectance and transmittance decreased exponentially as the transport optical thickness increased from 5.5 to 8.0. Unlike transmittance which decreased from 45% to 3%, reflectance was less sensitive to transport optical thickness and decreased by only 7.2%. Moreover, the exponential behaviors of transmittance and reflectance confirmed that transport optical thickness can be used as a reliable tool for determining laser propagation through the skin. It means by knowing the transport optical thickness at a given wavelength, transmittance and reflectance can be calculated using the correlations given in Figure 3.6.

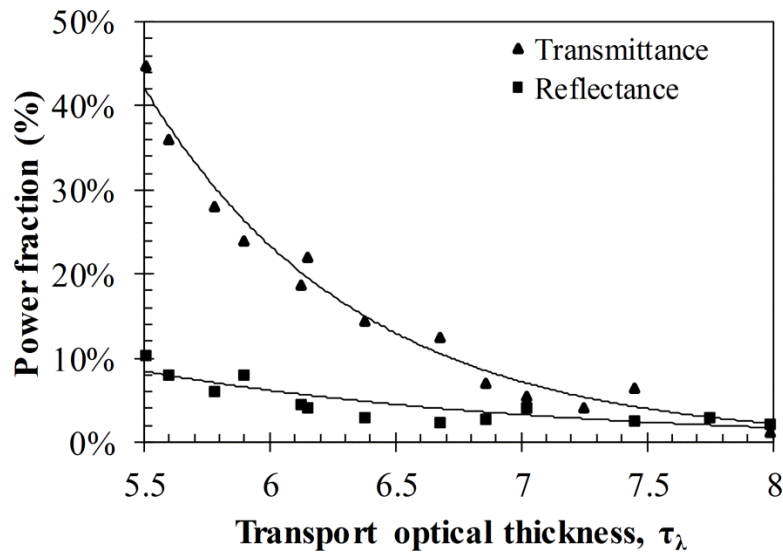


Figure 3.6: Correlation between transport optical thickness, transmittance and reflectance

Penetration depth is defined as the depth within the tissue at which the irradiance of laser falls to lower than 37% ( $1/e$ ) of the initial irradiance. This depth indicates how deep the laser energy is deposited within the tissue. In this study, in order to calculate the penetration depth forward Monte Carlo method was performed, using the recovered optical properties. Then, based on the density of the absorbed energy at each location, penetration depth was calculated. Figure 3.7 demonstrates the penetration depth of the human skin in the wavelength range from 1000 nm to 1900 nm. Based on this figure, the penetration depth reached its lowest value at the wavelength of 1480 nm. It means at this wavelength, absorption in layers of epidermis and dermis was dominant compared to subcutaneous. Moreover, for wavelengths higher than 1700 nm, the penetration depth was between 2 mm and 2.5 mm, which confirm the fact that subcutaneous is highly absorbent at this spectral range and laser energy is mostly absorbed at the very beginning

of this layer. On the other hand, in wavelength range from 1000 nm to 1450 nm, laser penetrated deeply through the human skin and the penetration depth was higher than 4.0 mm within this range. Finally, as it is shown in Figure 3.7, a wide range of penetration depths can be reached by using NIR lasers which ensures the applicability of this spectral range. Any penetration depth between 1.5 mm to 4.5 mm can be simply obtained by selecting the wavelength accordingly.

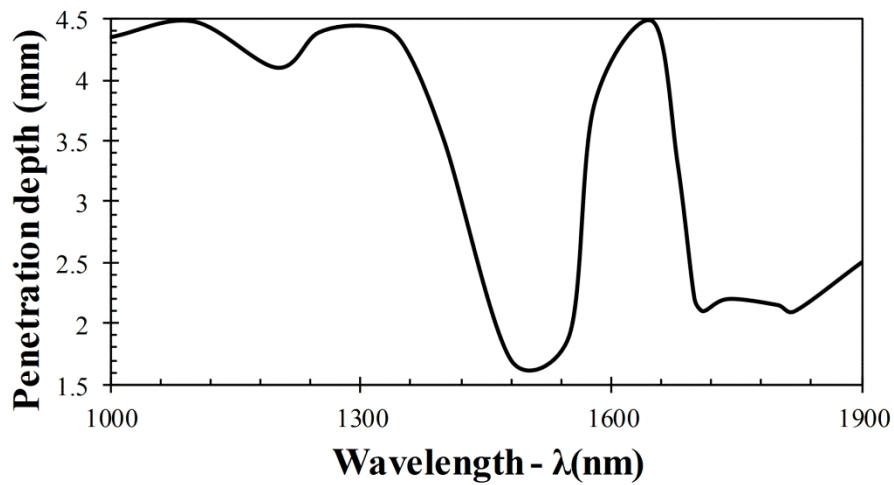


Figure 3.7: Penetration depth of laser propagation through human skin

### 3.2.3 Conclusion

In this chapter, a three-dimensional single-layer inverse Monte Carlo method was developed to calculate the optical properties of the skin using the experimental values of reflectance and transmittance. The implemented method was then validated against results of Prahl *et al.*, van de Hulst, Giovanelli and Bashkatov *et al.* For all these benchmark solutions the percent deviation was lower than 2%. Then, several sensitivity analyses were performed to ensure that assumption of constant asymmetry factor and

refractive index does not change the output results for more than 2%. Finally, by employing the reduced-variance forward Monte code in this inverse method, uncertainties of the results were minimized to 0.6%.

In the next step, by using this validated inverse method, the optical properties of human skin were recovered from the experimental results of Hardy *et al.* It was observed that both reflectance and transmittance scale well with transport optical thickness. The exponential correlation between these factors confirmed that transport optical thickness can be utilized as a scale for comparing laser propagation through the skin. The penetration depth in this wavelength range was then calculated by performing the forward Monte Carlo code for the recovered optical properties. The major findings of this chapter can be summarized as:

- Both transmittance and reflectance showed a dynamically increasing and decreasing behavior within the wavelength range from 1000 nm to 1900 nm.
- Unlike transmittance which decreased from 45% to 3%, reflectance was less sensitive to transport optical thickness and decreased by only 7.2%.
- Transport optical thickness scales well with both transmittance and reflectance and can be used as a tool to compare the laser penetration at different wavelengths.
- At wavelengths shorter than 1450 nm, laser penetrated deeply through the skin and the penetration depth was higher than 4.0 mm.

- At wavelengths higher than 1700 nm, the penetration depths were between 2 mm and 2.5 mm, which confirm the fact that subcutaneous is highly absorbent at this range.
- The shortest penetration depth can be achieved at the wavelength of 1480 nm where this depth fell to 1.7 mm. At this wavelength, epidermis and dermis absorbed the laser energy dominantly.

## Chapter 4

### Conclusion and Recommendations

#### 4.1 SUMMARY

This thesis presented numerical studies addressing the laser interaction with human skin. The study is mainly includes (i) laser propagation through the human skin using forward Monte Carlo method as well as (ii) wavelength selection analysis for medical purposes (i.e. LLLT) using inverse Monte Carlo method.

##### 4.1.1 Light propagation through the tissue

The first part of the study provided a three-dimensional multi-layer reduced-variance Monte Carlo method for simulating the laser propagation through the human skin in the wavelength range from 1000 nm to 1900 nm. Human skin was modeled as a three layer participating medium where its properties were obtained from literature. The major findings of this part of the study can be summarized as:

- Based on the selected wavelength, the laser power can be efficiently delivered to the specific layer of the skin. For example, 90% absorptance in subcutaneous and 48% absorptance in dermis can be obtained by using wavelengths of 1780 nm and 1550 nm, respectively.
- Although the total absorptance at the wavelengths of 1150 nm, 1350 nm and 1650 nm are fairly close, the spatial distribution of the absorbed power within the skin



layers is considerably different. To illustrate, dermis is the main absorber at the wavelength of 1350 nm, whereas in the cases of 1150 nm and 1650 nm, the incident power was mostly deposited to subcutaneous.

- Similar absorptance rate in one layer does not indicate the similar laser propagation. For example, absorptance level in dermis is fairly similar at both wavelengths of 1150 nm and 1650 nm. However, the volumetric dosage in dermis at wavelength of 1150 nm never exceeded the value of  $280 \text{ mW/mm}^3$ , while at wavelength of 1650 nm this value was as high as  $400 \text{ mW/mm}^3$ .
- In case of power delivery to dermis, wavelength of 1350 nm for dosages lower than  $150 \text{ mW/mm}^3$  and 1550 nm for dosages higher than  $250 \text{ mW/mm}^3$  should be utilized. In addition, wavelength of 1780 nm featured the highest power delivery to subcutaneous.
- At a given laser power, the local dosage of absorbed power can be increased about three times, by changing the beam profile from uniform (flat) to a sharp Gaussian's one.
- Maximum permissible exposure durations is inversely related to the laser power. For the laser power of 155 mW, the exposure durations can be up to 19 s and 15 s at wavelengths of 1150 nm and 1350 nm, respectively.

#### **4.1.2 Wavelength selection based on transmittance and reflectance**

In the second part of the study, a reduced-variance single-layer inverse Monte Carlo method was presented to retrieve the optical properties of human skin from the experimental values of transmittance and reflectance. The inverse method was then validated against the benchmark solutions for both isotropic and anisotropic cases. In addition, several sensitivity analyses were performed to minimize the error and uncertainty of the output results. Below are the major findings of this part of study:

- Both transmittance and reflectance showed a dynamically increasing and decreasing behavior within the wavelength range from 1000 nm to 1900 nm.
- Unlike transmittance which decreased from 45% to 3%, reflectance was less sensitive to transport optical thickness and decreased by only 7.2%.
- Transport optical thickness scales well with both transmittance and reflectance and can be used as a tool to compare the laser penetration at different wavelengths.
- At wavelengths shorter than 1450 nm, laser penetrated deeply through the skin and the penetration depth was higher than 4.0 mm.
- At wavelengths higher than 1700 nm, the penetration depths were between 2 mm and 2.5 mm, which confirm the fact that subcutaneous is highly absorbent at this range.
- The shortest penetration depth can be achieved at the wavelength of 1480 nm where this depth fell to 1.7 mm. At this wavelength, epidermis and dermis absorbed the laser energy dominantly.

## 4.2 RECOMMENDATIONS FOR FUTURE RESEARCH

The presented thesis opens several questions to be answered through further research. Following are the recommended studies that can be performed for further understanding of the discussed subject:

- The skin model used in this study was a three layer medium with constant optical properties at each layer. However, the optical properties of skin are gradually changed through the skin layers. Thus, a more accurate skin model with variable optical properties can provide more precise results.
- The muscle layer beneath the subcutaneous was considered to be a perfect absorber due to its high value of absorption coefficient. A more accurate simulation can be performed by considering the light penetration through muscle layer and deeper as well.
- Due to excessive increase in CPU time, the number of energy packets was restricted to  $10^6$  bundles. Therefore, more precise results can be achieved by utilizing supercomputers.
- The implemented inverse Monte Carlo method was capable of finding the optical properties for a single-layer skin. However, skin properties are not constant through the skin layers and they change considerably as function of depth. Thus, the next step of this study can be implementing a multi-layer inverse method which would be able to measure the optical properties of each layer separately.
- Asymmetry factor as well as refractive index varies as a function of wavelength, while they were considered constant in this study. Thus, for further studies, these

two factors can be assumed as unknown values besides absorption and scattering coefficients.

- As a future research, this presented study can be coupled to the experimental set-up to obtain the optical properties of an unknown medium. Since the implemented code can be extended to the multi-layer tool, combinations of different media can be studied as well.

## **Appendices**

## APPENDIX A: MATLAB SCRIPT FOR ORIGINAL MONTE CARLO METHOD

Following is the main code for forward Monte Carlo method used in this study. The inputs of this code can be listed as: number of layers, slab dimensions, scattering coefficient, absorption coefficient, asymmetry factor, refractive index, phase function and number of energy packets. This code generates the location of the absorbed packets within and beneath the skin.

```
1  %set number of layers
2  n_layer=5;
3
4  % set number of packets
5  N_total_packets=10^6;
6
7  % set the matrix of location
8  loc_abs=zeros(N_total_packets,3);
9
10 % Inputs
11 scattering_coeff=[0,141.2,141.2,94.8,0]; %cm-1
12 absorption_coeff=[1000,2.7,8.58,1.69,1000];%cm-1
13 geometry=[4,4,2;4,4,0.03;4,4,0.12;4,4,0.3;4,4,2];%cm
14 ref_index=[1,1.37,1.37,1.37,1.37];
15 %load optical properties of layers;
16 load('phase_function_skin')
17
18 for i=1:n_layer
19     beta(i)=scattering_coeff(i)+absorption_coeff(i);
20     albedo(i)=scattering_coeff(i)./beta(i);
21     length(i)=geometry(i,1);
22     width(i)=geometry(i,2);
23     thickness(i)=geometry(i,3);
24     n(i)=ref_index(i);
25 end
26
27 % matrix of inputs
28 input=zeros(n_layer,6);
29 input(:,1)=beta(1,:);
30 input(:,2)=albedo(1,:);
31 input(:,3)=length(1,:);
32 input(:,4)=width(1,:);
33 input(:,5)=thickness(1,:);
34 input(:,6)=n(1,:);
```

```

35
36
37 %set number of packet hits on each wall to zero
38 s_abs_bottom=0;
39 s_packet_exit_col_top=0;
40 s_packet_exit_noncol_top=0;
41 s_missed=0;
42 s_abs_within=0;
43 for ii=1:n_layer
44     str=num2str(ii);
45     eval(['s_abs_within_layer_' str '=0;']) ;
46 end
47
48 %set the start point at layer#2
49 i=2;
50
51 n_count=1;
52 while n_count<N_total_packets
53
54     % set initial position of packet
55     radius_initial=0.306*rand;
56     Angle_initial=2*pi*rand;
57     x=(0.5*length(2))+radius_initial*cos(Angle_initial);
58     y=(0.5*width(2))+radius_initial*sin(Angle_initial);
59     z=0;
60     %collimated radiation
61     theta=0;
62     phi=2*pi*rand;
63     %run until the packet gets absotbed or exits
64     packet_exit=0;
65     while packet_exit==0
66         packet_exit_layer=0;
67         while packet_exit_layer == 0;
68             if theta==0
69                 d_wall=thickness(i)-z;
70                 zhit=thickness(i);
71                 xhit=x;
72                 yhit=y;
73             elseif theta==pi
74                 d_wall=z;
75                 zhit=0;
76                 xhit=x;
77                 yhit=y;
78             else
79
80 [xhit,yhit,zhit,d_wall]=find_angles(x,y,z,phi,theta,length(i),
81 width(i),thickness(i));
82         end
83         path_length=(-log(rand))/beta(i);
84         if path_length>d_wall
85             if zhit==thickness(i) || zhit==0
86                 if theta==0

```

```

87         if i==n_layer-1
88             if
89 rand>reflectance(theta,n(i),n(i+1))
90                 s_abs_bottom=1+s_abs_bottom;
91                 loc_abs(n_count,1)=xhit;
92                 loc_abs(n_count,2)=yhit;
93
94 loc_abs(n_count,3)=zhit+thickness(2)+thickness(3);
95                 n_count=n_count+1;
96                 packet_exit=1;
97                 packet_exit_layer=1;
98             else
99                 x=x;
100                y=y;
101                z=thickness(i);
102                phi=phi;
103                theta=pi;
104            end
105        else
106            if
107 rand>reflectance(theta,n(i),n(i+1))
108                x=x;
109                y=y;
110                z=0;
111                theta=0;
112                phi=phi;
113                i=i+1;
114                packet_exit_layer=1;
115            else
116                x=x;
117                y=y;
118                z=thickness(i);
119                phi=phi;
120                theta=pi;
121            end
122        end
123
124        elseif theta==pi
125            if i==2
126                if
127 rand>reflectance(theta,n(i),n(i-1))
128
129 s_packet_exit_col_top=1+s_packet_exit_col_top;
130                 n_count=n_count+1;
131                 packet_exit=1;
132                 packet_exit_layer=1;
133             else
134                 x=x;
135                 y=y;
136                 z=0;
137                 phi=phi;
138                 theta=pi;

```



```

139         end
140     else
141         if
142     rand>reflectance(theta,n(i),n(i-1))
143             x=x;
144             y=y;
145             z=0;
146             theta=0;
147             phi=phi;
148             i=i-1;
149             packet_exit_layer=1;
150         else
151             x=x;
152             y=y;
153             z=0;
154             phi=phi;
155             theta=pi;
156         end
157     end
158
159     else
160         if zhit==thickness(i)
161             if i==n_layer-1
162                 if n(i)>n(i+1)
163                     if
164     theta>find_TIR(n(i),n(i+1));
165                         x=xhit;
166                         y=yhit;
167                         z=zhit;
168                         phi=phi;
169                         theta=pi-theta;
170                     else
171                         if
172     rand>reflectance(theta,n(i),n(i+1))
173
174     s_abs_bottom=s_abs_bottom+1;
175
176     loc_abs(n_count,1)=xhit;
177
178     loc_abs(n_count,2)=yhit;
179
180     loc_abs(n_count,3)=zhit+thickness(2)+thickness(3);
181                                     n_count=n_count+1;
182                                     packet_exit=1;
183
184     packet_exit_layer=1;
185         else
186             x=xhit;
187             y=yhit;
188             z=zhit;
189             phi=phi;
190             theta=pi-theta;

```

```

191                                     end
192                                 end
193                             else
194                                 if
195 rand>reflectance(theta,n(i),n(i+1))
196
197 s_abs_bottom=s_abs_bottom+1;
198
199 loc_abs(n_count,1)=xhit;
200
201 loc_abs(n_count,2)=yhit;
202
203 loc_abs(n_count,3)=zhit+thickness(2)+thickness(3);
204                                     n_count=n_count+1;
205                                     packet_exit=1;
206                                     packet_exit_layer=1;
207                                 else
208                                     x=xhit;
209                                     y=yhit;
210                                     z=zhit;
211                                     phi=phi;
212                                     theta=pi-theta;
213                                 end
214
215                             end
216                         else
217
218
219                             if n(i)>n(i+1)
220                                 if
221 theta>find_TIR(n(i),n(i+1));
222                                     x=xhit;
223                                     y=yhit;
224                                     z=zhit;
225                                     phi=phi;
226                                     theta=pi-theta;
227                                 else
228                                     if
229 rand>reflectance(theta,n(i),n(i+1))
230                                     x=x;
231                                     y=y;
232                                     z=0;
233                                     phi=phi;
234
235 theta=asin((n(i)/n(i+1))*sin(theta));
236                                     i=i+1;
237
238 packet_exit_layer=1;
239
240                                 else
241                                     x=xhit;
242                                     y=yhit;

```

```

243                                     z=zhit;
244                                     phi=phi;
245                                     theta=pi-theta;
246                                     end
247                                     end
248                                 else
249                                     if
250 rand>reflectance(theta,n(i),n(i+1))
251                                     x=x;
252                                     y=y;
253                                     z=0;
254                                     phi=phi;
255
256 theta=asin((n(i)/n(i+1))*sin(theta));
257                                     i=i+1;
258                                     packet_exit_layer=1;
259                                 else
260                                     x=xhit;
261                                     y=yhit;
262                                     z=zhit;
263                                     phi=phi;
264                                     theta=pi-theta;
265                                 end
266
267                                     end
268                                 end
269
270
271                                 elseif zhit==0
272                                     if i==2
273                                         if n(i)>n(i-1)
274                                             if pi-
275 theta>find_TIR(n(i),n(i-1));
276                                             x=xhit;
277                                             y=yhit;
278                                             z=zhit;
279                                             phi=phi;
280                                             theta=pi-theta;
281                                         else
282                                             if
283 rand>reflectance(theta,n(i),n(i-1))
284
285 s_packet_exit_noncol_top=s_packet_exit_noncol_top+1;
286                                             n_count=n_count+1;
287                                             packet_exit=1;
288
289 packet_exit_layer=1;
290                                         else
291                                             x=xhit;
292                                             y=yhit;
293                                             z=zhit;
294                                             phi=phi;

```

```

295                                     theta=pi-theta;
296                                     end
297                                 end
298                             else
299                                 if
300 rand>reflectance(theta,n(i),n(i-1))
301
302 s_packet_exit_noncol_top=s_packet_exit_noncol_top+1;
303                                     n_count=n_count+1;
304                                     packet_exit=1;
305                                     packet_exit_layer=1;
306                                 else
307                                     x=xhit;
308                                     y=yhit;
309                                     z=zhit;
310                                     phi=phi;
311                                     theta=pi-theta;
312                                 end
313
314                             end
315                         else
316
317
318                                 if n(i)>n(i+1)
319                                     if pi-
320 theta>find_TIR(n(i),n(i-1));
321                                     x=xhit;
322                                     y=yhit;
323                                     z=zhit;
324                                     phi=phi;
325                                     theta=pi-theta;
326                                 else
327                                     if
328 rand>reflectance(theta,n(i),n(i-1))
329                                     x=x;
330                                     y=y;
331                                     z=thickness(i-1);
332                                     phi=phi;
333                                     theta=pi-
334 asin((n(i)/n(i-1))*sin(pi-theta));
335                                     i=i-1;
336
337 packet_exit_layer=1;
338                                 else
339                                     x=xhit;
340                                     y=yhit;
341                                     z=zhit;
342                                     phi=phi;
343                                     theta=pi-theta;
344                                 end
345                             end
346                         else

```

```

347                                     if
348 rand>reflectance(theta,n(i),n(i-1))
349                                     x=x;
350                                     y=y;
351                                     z=0;
352                                     phi=phi;
353                                     theta=pi-
354 asin((n(i)/n(i-1))*sin(pi-theta));
355                                     i=i-1;
356                                     packet_exit_layer=1;
357                                     else
358                                     x=xhit;
359                                     y=yhit;
360                                     z=zhit;
361                                     phi=phi;
362                                     theta=pi-theta;
363                                     end
364                                     end
365                                     end
366                                     end
367                                     end
368
369                                     end
370                                     elseif xhit==0
371                                     s_missed=s_missed+1;
372                                     n_count=n_count+1;
373                                     packet_exit=1;
374                                     packet_exit_layer=1;
375                                     elseif xhit==length(i)
376                                     s_missed=s_missed+1;
377                                     n_count=n_count+1;
378                                     packet_exit=1;
379                                     packet_exit_layer=1;
380                                     elseif yhit==0
381                                     s_missed=s_missed+1;
382                                     n_count=n_count+1;
383                                     packet_exit=1;
384                                     packet_exit_layer=1;
385                                     elseif yhit==width(i)
386                                     s_missed=s_missed+1;
387                                     n_count=n_count+1;
388                                     packet_exit=1;
389                                     packet_exit_layer=1;
390                                     end
391                                     else
392                                     frac_traveled=path_length/d_wall;
393                                     if theta==0
394                                     x=x;
395                                     y=y;
396                                     z=z+path_length;
397                                     elseif theta==pi

```

```

399         x=x;
400         y=y;
401         z=z-path_length;
402     else
403         x=x+(xhit-x)*frac_traveled;
404         y=y+(yhit-y)*frac_traveled;
405         z=z+(zhit-z)*frac_traveled;
406     end
407     if rand<albedo(i)
408         R_sca=rand;
409         kk=1;
410         while R_sca>=phase_function_skin(kk,7)
411
412 scattering_angle=(phase_function_skin(kk,3)+phase_function_ski
413 n(kk+1,3))/2;
414             kk=kk+1;
415         end
416         phi_2=2*pi*rand;
417         theta=theta+scattering_angle*cos(phi_2);
418         phi=phi+scattering_angle*sin(phi_2);
419     else
420         ii=i;
421         str=num2str(ii);
422         eval(['s_abs_within_layer_' str
423 '=1+s_abs_within_layer_' str ';' ]) ;
424         loc_abs(n_count,1)=x;
425         loc_abs(n_count,2)=y;
426         if i==2
427             loc_abs(n_count,3)=z;
428         else
429             ooo=0;
430             for u=2:i-1
431                 ooo=ooo+thickness(u);
432             end
433             loc_abs(n_count,3)=z+ooo;
434         end
435         s_abs_within=s_abs_within+1;
436         n_count=n_count+1;
437         packet_exit=1;
438         packet_exit_layer=1;
439     end
440 end
441 end
442
443 end
444 end
445
446 % Tally up the absorped packets(within and bottom), the missed
447 packets, exited from
448 % top(collimated and diffuese)

```

```
449 check_total_packets=  
450 s_abs_bottom+s_packet_exit_col_top+s_packet_exit_noncol_top+s_  
451 missed+s_abs_within;  
452 absorptance=(s_abs_bottom+s_abs_within)/N_total_packets;  
453 transmitted=s_abs_bottom/N_total_packets;
```

## APPENDIX B: MATLAB SCRIPT FOR REDUCED-VARIANCE MONTE CARLO METHOD

Following is the code for reduced-variance Monte Carlo method used in this study. The inputs of this code can be listed as: number of layers, slab dimensions, scattering coefficient, absorption coefficient, asymmetry factor, refractive index, phase function, number of energy packets and threshold energy. This code generates the location of the absorbed packets within and beneath the skin.

```
1  clc
2  clear
3
4  %set number of layers,in this case the first and the last
5  layers are
6  %considered to be ambient (air)
7  n_layer=5;
8
9  % set number of packets
10 N_total_packets=10^6;
11
12 % set the matrix of location
13 loc_abs=zeros(N_total_packets,3);
14
15 % Inputs
16 absorption_coeff=[1000,1.5,6.4,6.1,1000];%cm-1
17 scattering_coeff=[0,132.9,132.9,85.9,0]; %cm-1
18 geometry=[4,4,2;4,4,0.03;4,4,0.12;4,4,0.3;4,4,2];%cm
19 ref_index=[1,1.34,1.4,1.44,1.44];
20 %load optical properties of layers;
21 load('phase_function_skin')
22
23 for i=1:n_layer
24     beta(i)=scattering_coeff(i)+absorption_coeff(i);
25     albedo(i)=scattering_coeff(i)./beta(i);
26     length(i)=geometry(i,1);
27     width(i)=geometry(i,2);
28     thickness(i)=geometry(i,3);
29     n(i)=ref_index(i);
30 end
31
32 % matrix of inputs
33 input=zeros(n_layer,6);
34 input(:,1)=beta(1,:);
35 input(:,2)=albedo(1,:);
36 input(:,3)=length(1,:);
37 input(:,4)=width(1,:);
```



```

38 input(:,5)=thickness(1,:);
39 input(:,6)=n(1,:);
40
41
42 %set number of packet hits on each wall to zero
43 w_abs_bottom=0;
44 w_packet_exit_col_top=0;
45 w_packet_exit_noncol_top=0;
46 w_missed=0;
47 w_abs_within=0;
48 w_threshold=0.001;
49
50 %set the start point at layer#2, considering i=1 is ambient
51 i=2;
52
53 n_count=1;
54 while n_count<N_total_packets
55     while w>w_threshold && m<=10
56         m=0; % Number of Interactions
57         w=1; % Weight of the packets
58         % set initial position of packet
59         radius_initial=0.306*rand;
60         Angle_initial=2*pi*rand;
61         x=(0.5*length(2))+radius_initial*cos(Angle_initial);
62         y=(0.5*width(2))+radius_initial*sin(Angle_initial);
63         z=0;
64         %collimated radiation
65         theta=0;
66         phi=2*pi*rand;
67         %run until the packet gets absotbed or exits
68         packet_exit=0;
69         while packet_exit==0
70             packet_exit_layer=0;
71             while packet_exit_layer == 0;
72                 % Measuring the distance from the wall and
73 also the
74                 % potential hittin point
75                 if theta==0
76                     d_wall=thickness(i)-z;
77                     zhit=thickness(i);
78                     xhit=x;
79                     yhit=y;
80                 elseif theta==pi
81                     d_wall=z;
82                     zhit=0;
83                     xhit=x;
84                     yhit=y;
85                 else
86
87 [xhit,yhit,zhit,d_wall]=find_angles(x,y,z,phi,theta,length(i),
88 width(i),thickness(i));
89                 end

```

```

90                                     % Calculate the path length to find out if the
91 packet will
92                                     % hit the wall or not
93 path_length=(-log(rand))/beta(i);
94 if path_length>d_wall
95     if zhit==thickness(i) || zhit==0
96         if theta==0 % If a collimated packet
97 hits the bottom or top surface of a layer
98         if i==n_layer-1 % hitting hte last
99 layer above ambient
100                                     if
101 rand>reflectance(theta,n(i),n(i+1))
102
103 w_abs_bottom=w+w_abs_bottom;
104                                     loc_abs(n_count,1)=xhit;
105                                     loc_abs(n_count,2)=yhit;
106
107 loc_abs(n_count,3)=zhit+thickness(2)+thickness(3);
108                                     m=m+1;
109                                     n_count=n_count+1;
110                                     packet_exit=1;
111                                     packet_exit_layer=1;
112     else
113
114 w_abs_bottom=w_abs_bottom+(1-albedo(i))*w;
115                                     w=albedo(i)*w;
116                                     m=m+1;
117                                     x=x;
118                                     y=y;
119                                     z=thickness(i);
120                                     phi=phi;
121                                     theta=pi;
122     end
123 else
124
125     if
126 rand>reflectance(theta,n(i),n(i+1))
127                                     x=x;
128                                     y=y;
129                                     z=0;
130                                     theta=0;
131                                     phi=phi;
132                                     i=i+1;
133                                     packet_exit_layer=1;
134     else
135
136                                     x=x;
137                                     y=y;
138                                     z=thickness(i);
139                                     phi=phi;
140                                     theta=pi;

```

```

141
142 w_abs_within=w_abs_within+(1-albedo(i))*w;
143                                     m=m+1;
144                                     end
145                                     end
146
147                                     elseif theta==pi
148                                         if i==2
149                                             if
150 rand>reflectance(theta,n(i),n(i-1))
151
152 w_packet_exit_col_top=w+w_packet_exit_col_top;
153                                     m=m+1;
154                                     n_count=n_count+1;
155                                     packet_exit=1;
156                                     packet_exit_layer=1;
157                                     else
158                                         x=x;
159                                         y=y;
160                                         z=0;
161                                         phi=phi;
162                                         theta=0;
163
164                                     end
165                                     else
166                                         if
167 rand>reflectance(theta,n(i),n(i-1))
168                                         x=x;
169                                         y=y;
170                                         z=0;
171                                         theta=pi;
172                                         phi=phi;
173                                         i=i-1;
174                                         packet_exit_layer=1;
175                                     else
176                                         x=x;
177                                         y=y;
178                                         z=0;
179                                         phi=phi;
180                                         theta=0;
181                                     end
182                                     end
183
184                                     else
185                                         if zhit==thickness(i)
186                                             if i==n_layer-1
187                                                 if n(i)>n(i+1)
188                                                     if
189 theta>find_TIR(n(i),n(i+1));
190                                                         x=xhit;
191                                                         y=yhit;
192                                                         z=zhit;

```

```

193                                     phi=phi;
194                                     theta=pi-theta;
195
196 w_abs_within=w_abs_within+(1-albedo(i))*w;
197                                     w=albedo(i)*w;
198                                     m=m+1;
199                                     else
200                                     if
201 rand>reflectance(theta,n(i),n(i+1))
202
203 w_abs_bottom=w_abs_bottom+w;
204                                     m=m+1;
205
206 loc_abs(n_count,1)=xhit;
207
208 loc_abs(n_count,2)=yhit;
209
210 loc_abs(n_count,3)=zhit+thickness(2)+thickness(3);
211
212 n_count=n_count+1;
213                                     packet_exit=1;
214
215 packet_exit_layer=1;
216                                     else
217                                     x=xhit;
218                                     y=yhit;
219                                     z=zhit;
220                                     phi=phi;
221                                     theta=pi-
222 theta;
223
224 w_abs_bottom=w_abs_bottom+(1-albedo(i))*w;
225                                     w=w*albedo(i);
226                                     m=m+1;
227                                     end
228                                     end
229                                     else
230                                     if
231 rand>reflectance(theta,n(i),n(i+1))
232
233 w_abs_bottom=w_abs_bottom+w;
234                                     m=m+1;
235
236 loc_abs(n_count,1)=xhit;
237
238 loc_abs(n_count,2)=yhit;
239
240 loc_abs(n_count,3)=zhit+thickness(2)+thickness(3);
241                                     n_count=n_count+1;
242                                     packet_exit=1;
243
244 packet_exit_layer=1;
245                                     else

```

```

246                                     x=xhit;
247                                     y=yhit;
248                                     z=zhit;
249                                     phi=phi;
250                                     theta=pi-theta;
251
252 w_abs_bottom=w_abs_bottom+(1-albedo(i))*w;
253                                     w=w*albedo(i);
254                                     m=m+1;
255                                     end
256
257                                     end
258 else
259
260                                     if n(i)>n(i+1)
261                                     if
262 theta>find_TIR(n(i),n(i+1));
263                                     x=xhit;
264                                     y=yhit;
265                                     z=zhit;
266                                     phi=phi;
267                                     theta=pi-theta;
268                                     else
269                                     if
270 rand>reflectance(theta,n(i),n(i+1))
271                                     x=x;
272                                     y=y;
273                                     z=0;
274                                     phi=phi;
275
276 theta=asin((n(i)/n(i+1))*sin(theta));
277                                     i=i+1;
278
279 w_abs_within=w_abs_within+(1-albedo(i))*w;
280                                     w=albedo(i)*w;
281                                     m=m+1;
282
283 packet_exit_layer=1;
284
285
286                                     else
287                                     x=xhit;
288                                     y=yhit;
289                                     z=zhit;
290                                     phi=phi;
291                                     theta=pi-
292 theta;
293
294 w_abs_within=w_abs_within+(1-albedo(i))*w;
295                                     w=albedo(i)*w;
296                                     m=m+1;
297                                     end

```

```

298                                     end
299                                     else
300                                     if
301 rand>reflectance(theta,n(i),n(i+1))
302                                     x=x;
303                                     y=y;
304                                     z=0;
305                                     phi=phi;
306
307 theta=asin((n(i)/n(i+1))*sin(theta));
308                                     i=i+1;
309
310 w_abs_within=w_abs_within+(1-albedo(i))*w;
311                                     w=albedo(i)*w;
312
313 packet_exit_layer=1;
314                                     else
315                                     x=xhit;
316                                     y=yhit;
317                                     z=zhit;
318                                     phi=phi;
319                                     theta=pi-theta;
320
321 w_abs_within=w_abs_within+(1-albedo(i))*w;
322                                     w=albedo(i)*w;
323                                     m=m+1;
324                                     end
325
326                                     end
327                                     end
328
329                                     elseif zhit==0
330                                     if i==2
331                                     if n(i)>n(i-1)
332                                     if pi-
333 theta>find_TIR(n(i),n(i-1));
334                                     x=xhit;
335                                     y=yhit;
336                                     z=zhit;
337                                     phi=phi;
338                                     theta=pi-theta;
339
340 w_abs_within=w_abs_within+(1-albedo(i))*w;
341                                     w=albedo(i)*w;
342                                     m=m+1;
343                                     else
344                                     if
345 rand>reflectance(theta,n(i),n(i-1))
346
347 w_packet_exit_noncol_top=w_packet_exit_noncol_top+w;
348                                     m=m+1;

```

```

349
350 n_count=n_count+1;
351                                     packet_exit=1;
352
353 packet_exit_layer=1;
354                                     else
355                                     x=xhit;
356                                     y=yhit;
357                                     z=zhit;
358                                     phi=phi;
359                                     theta=pi-
360 theta;
361
362 w_abs_within=w_abs_within+(1-albedo(i))*w;
363                                     w=albedo(i)*w;
364                                     m=m+1;
365                                     end
366                                     end
367                                     else
368                                     if
369 rand>reflectance(theta,n(i),n(i-1))
370
371 w_packet_exit_noncol_top=w_packet_exit_noncol_top+1;
372                                     m=m+1;
373                                     n_count=n_count+1;
374                                     packet_exit=1;
375
376 packet_exit_layer=1;
377                                     else
378                                     x=xhit;
379                                     y=yhit;
380                                     z=zhit;
381                                     phi=phi;
382                                     theta=pi-theta;
383
384 w_abs_within=w_abs_within+(1-albedo(i))*w;
385                                     w=albedo(i)*w;
386                                     m=m+1;
387                                     end
388
389                                     end
390 else
391
392                                     if n(i)>n(i+1)
393                                     if pi-
394 theta>find_TIR(n(i),n(i-1));
395                                     x=xhit;
396                                     y=yhit;
397                                     z=zhit;
398                                     phi=phi;
399                                     theta=pi-theta;

```

```

400
401 w_abs_within=w_abs_within+(1-albedo(i))*w;
402
403                                     w=albedo(i)*w;
404                                     m=m+1;
405                                     else
406                                     if
407 rand>reflectance(theta,n(i),n(i-1))
408                                     x=x;
409                                     y=y;
410                                     z=thickness(i-
411 1);
412                                     phi=phi;
413                                     theta=pi-
414 asin((n(i)/n(i-1))*sin(pi-theta));
415 w_abs_within=w_abs_within+(1-albedo(i))*w;
416                                     w=albedo(i)*w;
417                                     m=m+1;
418                                     i=i-1;
419
420 packet_exit_layer=1;
421                                     else
422                                     x=xhit;
423                                     y=yhit;
424                                     z=zhit;
425                                     phi=phi;
426                                     theta=pi-
427 theta;
428
429 w_abs_within=w_abs_within+(1-albedo(i))*w;
430                                     w=albedo(i)*w;
431                                     m=m+1;
432                                     end
433                                     end
434                                     else
435                                     if
436 rand>reflectance(theta,n(i),n(i-1))
437                                     x=x;
438                                     y=y;
439                                     z=0;
440                                     phi=phi;
441                                     theta=pi-
442 asin((n(i)/n(i-1))*sin(pi-theta));
443
444 w_abs_within=w_abs_within+(1-albedo(i))*w;
445                                     w=albedo(i)*w;
446                                     m=m+1;
447                                     i=i-1;
448
449 packet_exit_layer=1;
450                                     else
451                                     x=xhit;
452                                     y=yhit;

```



```

453                                     z=zhit;
454                                     phi=phi;
455                                     theta=pi-theta;
456
457 w_abs_within=w_abs_within+(1-albedo(i))*w;
458                                     w=albedo(i)*w;
459                                     m=m+1;
460                                     end
461
462                                     end
463                                 end
464                             end
465
466
467                                 end
468                             elseif xhit==0
469                                 w_missed=w_missed+1;
470                                 n_count=n_count+1;
471                                 packet_exit=1;
472                                 packet_exit_layer=1;
473                             elseif xhit==length(i)
474                                 w_missed=w_missed+1;
475                                 n_count=n_count+1;
476                                 packet_exit=1;
477                                 packet_exit_layer=1;
478                             elseif yhit==0
479                                 w_missed=w_missed+1;
480                                 n_count=n_count+1;
481                                 packet_exit=1;
482                                 packet_exit_layer=1;
483                             elseif yhit==width(i)
484                                 w_missed=w_missed+1;
485                                 n_count=n_count+1;
486                                 packet_exit=1;
487                                 packet_exit_layer=1;
488                             end
489                         else
490                             frac_traveled=path_length/d_wall;
491                             if theta==0
492                                 x=x;
493                                 y=y;
494                                 z=z+path_length;
495                             elseif theta==pi
496                                 x=x;
497                                 y=y;
498                                 z=z-path_length;
499                             else
500                                 x=x+(xhit-x)*frac_traveled;
501                                 y=y+(yhit-y)*frac_traveled;
502                                 z=z+(zhit-z)*frac_traveled;
503                             end
504

```

```

505         R_sca=rand;
506         kk=1;
507         while R_sca>=phase_function_skin(kk,7*(i-
508 1))
509
510 scattering_angle=(phase_function_skin(kk,(3+7*(i-
511 2)))+phase_function_skin(kk+1,(3+7*(i-2))))/2;
512         kk=kk+1;
513         end
514         phi_2=2*pi*rand;
515         theta=theta+scattering_angle*cos(phi_2);
516         phi=phi+scattering_angle*sin(phi_2);
517         w_abs_within=w_abs_within+(1-albedo(i))*w;
518         w=albedo(i)*w;
519         m=m+1;
520
521         end
522     end
523
524     end
525 end
526 [w,m]=russian_roulette(w);
527 end
528
529 % Tally up the absorped packets(within and bottom), the missed
530 packets, exited from
531 % top(collimated and diffuese)
532 check_total_packets=
533 w_abs_bottom+w_packet_exit_col_top+s_packet_exit_noncol_top+w_
534 missed+s_abs_within;
535 absorptance=(w_abs_bottom+s_abs_within)/N_total_packets;
536 transmitted=w_abs_bottom/N_total_packets;

```

## APPENDIX C: MATLAB SCRIPT FOR INVERSE MONTE CARLO METHOD

Following is the code for recovering the optical properties of human skin. The inputs of this code can be listed as: Transmittance and Reflectance.

```
1 clear
2 clc
3 % Input Parameters
4 T_input=0.4;
5 R_input=0.05;
6 emax=0.05;
7 % Set the initial guesses using Lookup Table
8 [kappa,sigma]= look_up_table(T,R);
9
10 while e<=emax
11 [T,R]=Monte_Carlo_One_Layer(kappa,sigma);
12 e=(((T-T_input)/T)^2 + ((R-R_input)/R)^2)^0.5;
13 [T,R]= interior_update(T,R,e);
14 end
15 T_final=T;
16 R_final=R;
```

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## **Vita**

Babak Nasouri received a Bachelor of Science in Mechanical Engineering from Sharif University of Technology in June 2012. His bachelor thesis was on simulating and characterizing the cycle variability of Homogeneous Charge Compression Ignition (HCCI) combustion. He began graduate study in the Thermal/Fluid Systems track of Mechanical Engineering department at The University of Texas at Austin in the fall of 2012. In January 2013, he began his research with Dr. Halil Berberoglu in the Solar Energy and Biofuels Laboratory, where his study was focused on laser propagation through participating media.

Permanent Email Address: [babak.nasouri@gmail.com](mailto:babak.nasouri@gmail.com)

This thesis was typed by the author.